

# The Radiative Return and Form Factors at Large $Q^2$

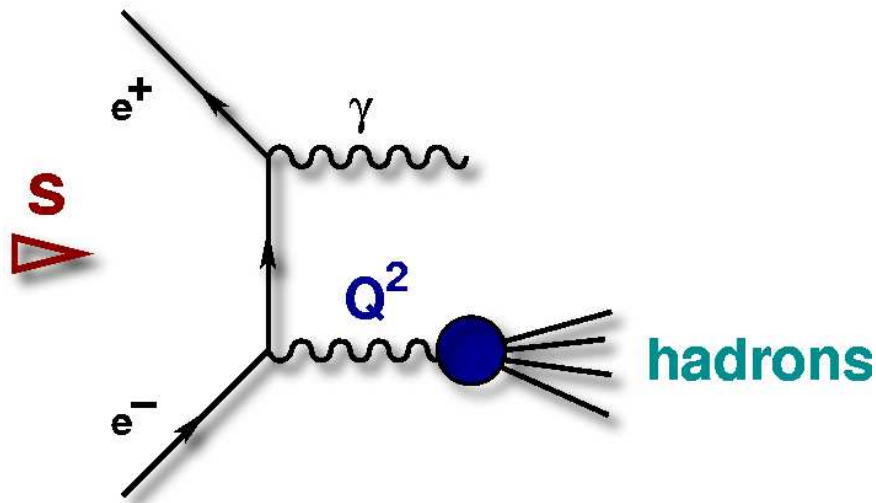
J.H. KÜHN, TTP, KARLSRUHE

- I Basic Idea
- II Monte Carlo Generators: Status & Perspectives
- III Nucleon Form Factors at B-Factories
- IV Pion and Kaon Form Factors at large  $Q^2$   
and  $\tau \rightarrow \nu K^- K^0$
- V Conclusions

# I BASIC IDEA

photon radiated off the initial  $e^+e^-$  (ISR) reduces the effective energy of the collision

$$d\sigma(e^+e^- \rightarrow \text{hadrons} + \gamma) = H(Q^2, \theta_\gamma) d\sigma(e^+e^- \rightarrow \text{hadrons})$$



- ▶ measurement of  $R(s)$  over the full range of energies, from threshold up to  $\sqrt{s}$
- ▶ large luminosities of factories compensate  $\alpha/\pi$  from photon radiation
- ▶ radiative corrections essential (NLO)
- ▶ advantage over energy scan (BES, CMD2, SND): systematics (e.g. normalization) only once

**High precision measurement of the hadronic cross-section at DAΦNE, CLEO-C, B-factories**

# Rough estimates for rates:

$\pi^+ \pi^- \gamma : E_\gamma > 100 \text{ MeV}$

$\sqrt{s}$ [GeV]	$\int \mathcal{L}$ [ $\text{fb}^{-1}$ ]	#events, $\theta_{min} = 7^\circ$
1.02	1.35	$16 \cdot 10^6$
10.6	100	$3.5 \cdot 10^6$

multi-hadron-events ( $R \equiv 2$ )  $\sqrt{s} = 10.6 \text{ GeV}$

$Q^2$ -interval [GeV]	#events, $\theta_{min} = 7^\circ$
[ 1.5 , 2.0 ]	$9.9 \cdot 10^5$
[ 2.0 , 2.5 ]	$7.9 \cdot 10^5$
[ 2.5 , 3.0 ]	$6.6 \cdot 10^5$
[ 3.0 , 3.5 ]	$5.8 \cdot 10^5$

# Lowest order

$$\frac{d\sigma}{dQ^2} (e^+e^- \rightarrow \gamma + \text{had}(Q^2)) = \sigma (e^+e^- \rightarrow \text{had}(Q^2))$$

$$\times \frac{\alpha}{\pi s} \left\{ \begin{array}{l} \frac{s^2+Q^4}{s(s-Q^2)} (\log(s/m_e^2) - 1) , \text{ no angular cut} \\ \frac{s^2+Q^4}{s(s-Q^2)} \log \left( \frac{1+\cos \theta_{min}}{1-\cos \theta_{min}} \right) - \frac{s-Q^2}{s} \cos \theta_{min} \end{array} \right\}$$

$$\Rightarrow \text{differential luminosity: } \frac{dL}{dQ^2} (Q^2, s) = \frac{\alpha}{\pi s} \left\{ \dots \right\} L(\text{at } s)$$

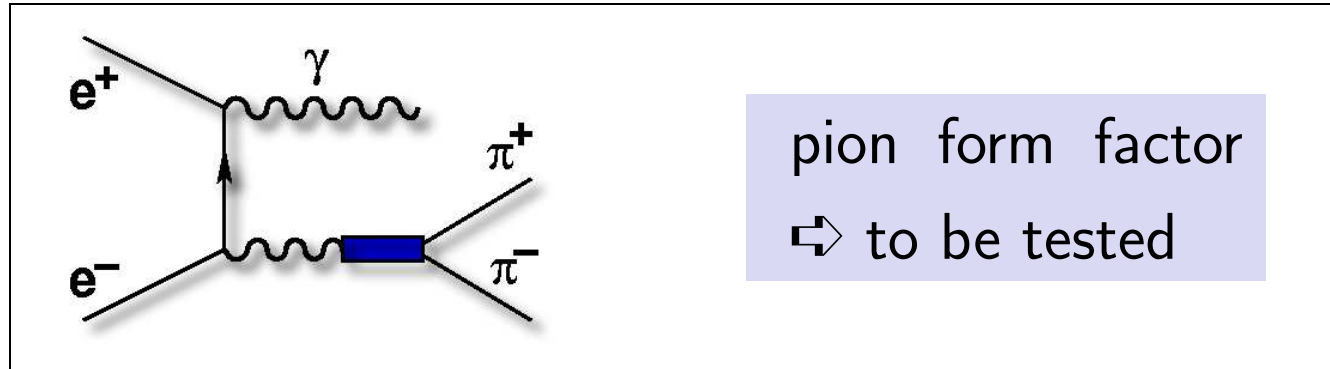
e.g.  $\theta_{min} = 30^\circ$  ;  $\sqrt{s} = 10.58$  GeV ;  $Q = 1$  GeV ;  $\Delta Q = 0.1$  GeV

$$\frac{dL}{dQ^2} (Q^2, s) \Delta Q^2 = 7.6 \cdot 10^{-6} L(\text{at } s)$$

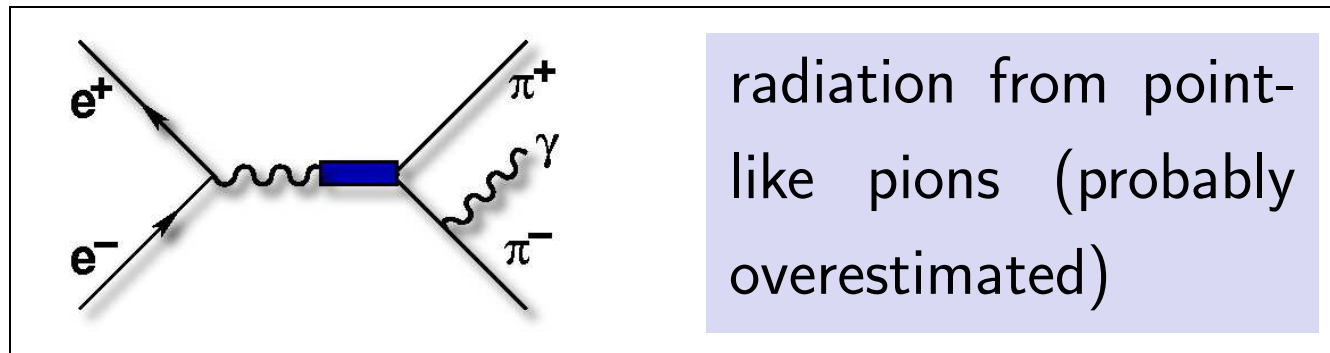
$100 \text{ fb}^{-1}$  at 10.58 GeV  $\Rightarrow$   $0.76 \text{ pb}^{-1}$  per scan point at 1 GeV

# Basic Ingredients for Pion Formfactor

## ► ISR

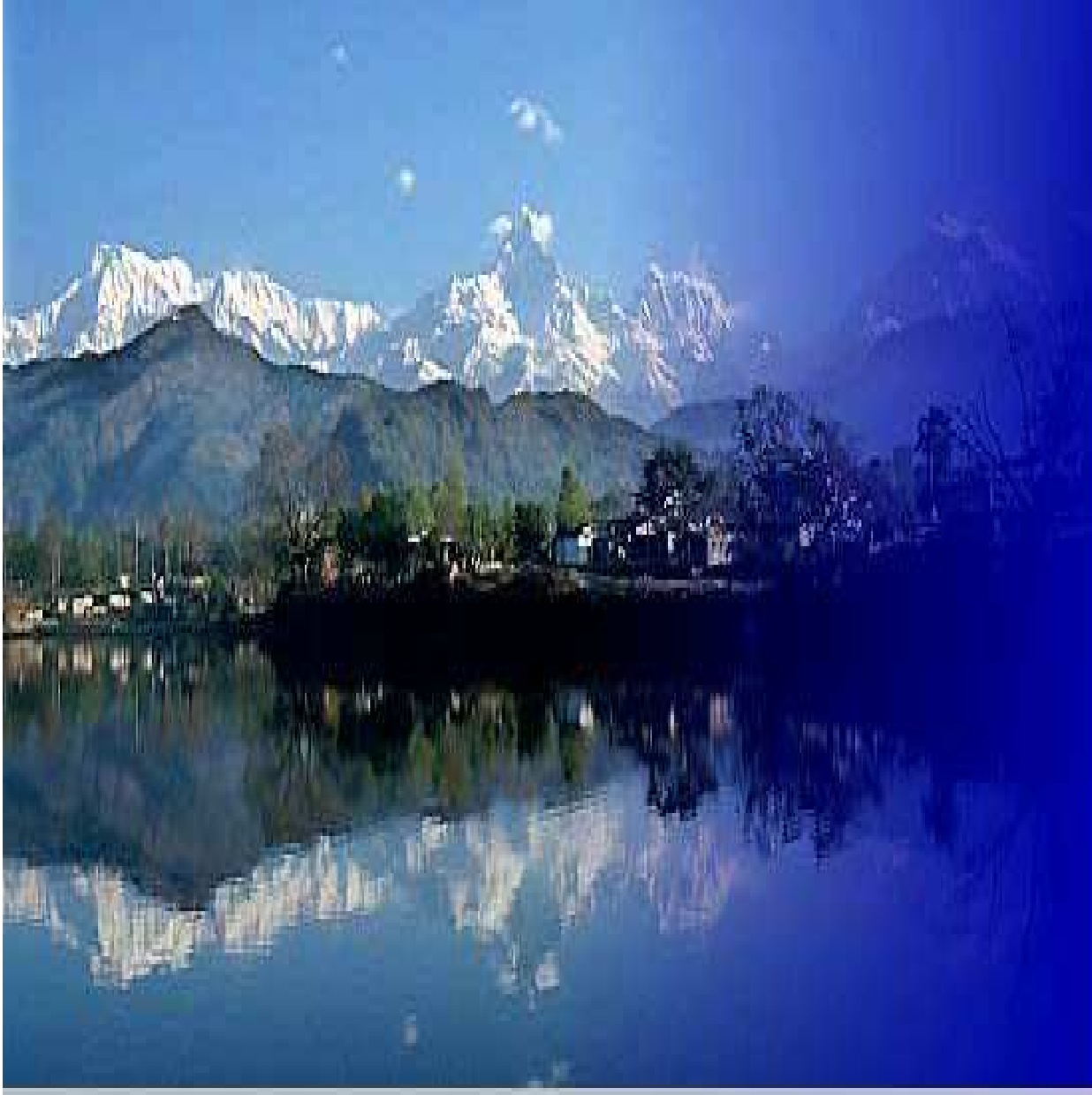


## ► FSR



- **additional radiation:** collinear (EVA MC)  
or NLO calculation (PHOKHARA MC)

# II MONTE CARLO GENERATORS



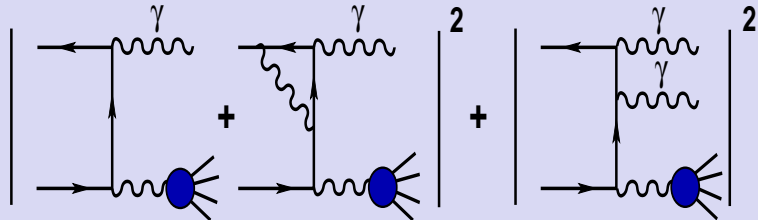
P  
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OTONS FROM  
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References etc. → <http://cern.ch/german.rodrico/phokhara>

## PHOKHARA 2.0:

$$\pi^+\pi^-, \mu^+\mu^-, 4\pi$$

- **ISR at NLO:** virtual corrections to one photon events and two photon emission at tree level

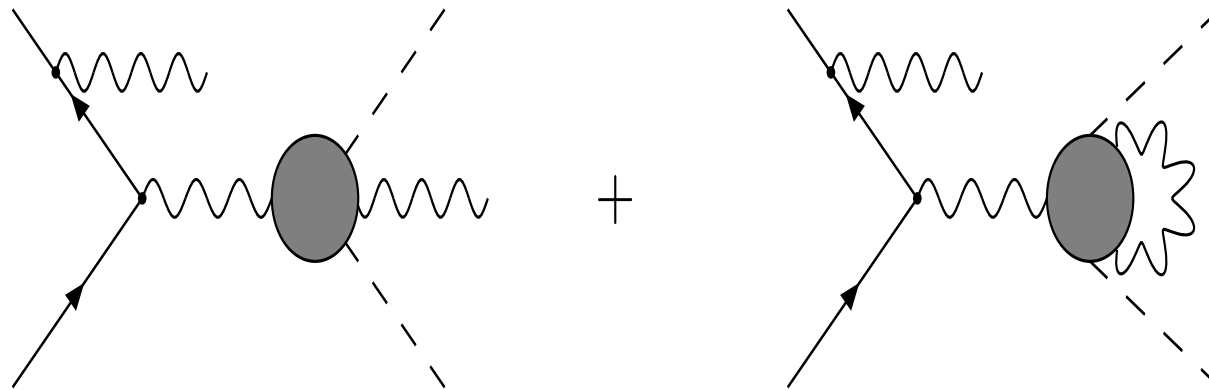


- FSR at LO:  $\pi^+\pi^-, \mu^+\mu^-$
- tagged or untagged photons
- modular structure

- ① LL at a fixed order + subleading terms (1 %)
- ② Full angular dependence
- ③ Momentum conservation
- ④ Tagged or untagged photon

# PHOKHARA 3.0

- ▶ specifically developed for  $\pi^+\pi^-$  (plus photons)
- ▶ allows for **simultaneous** emission of photons from **initial and final state**, including virtual corrections (interference neglected).

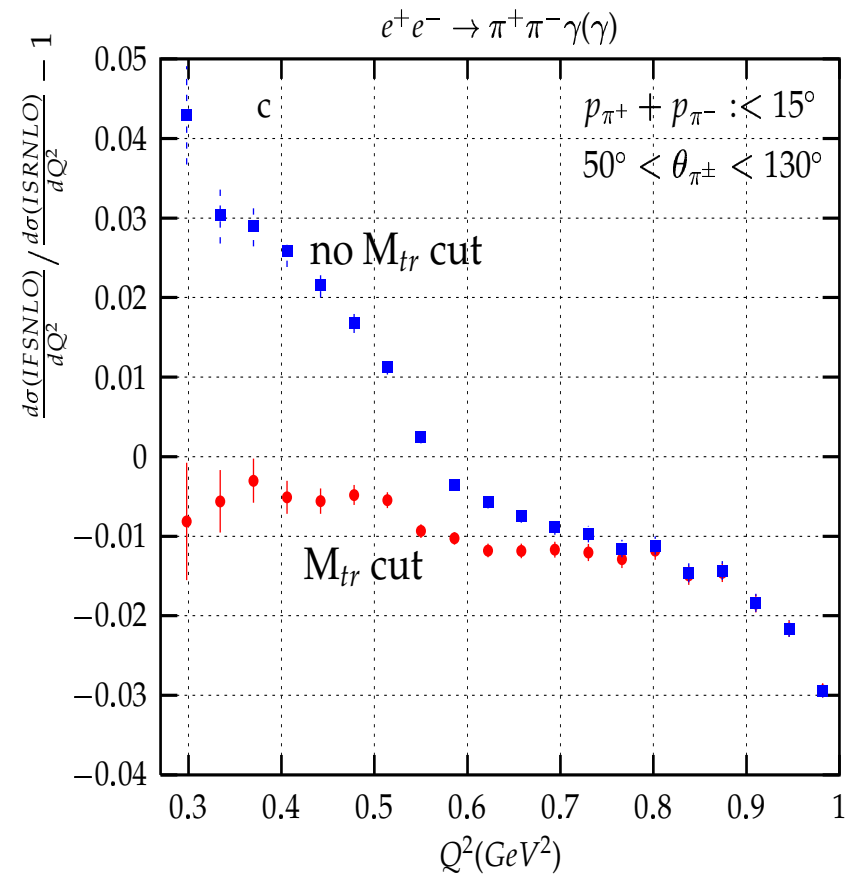
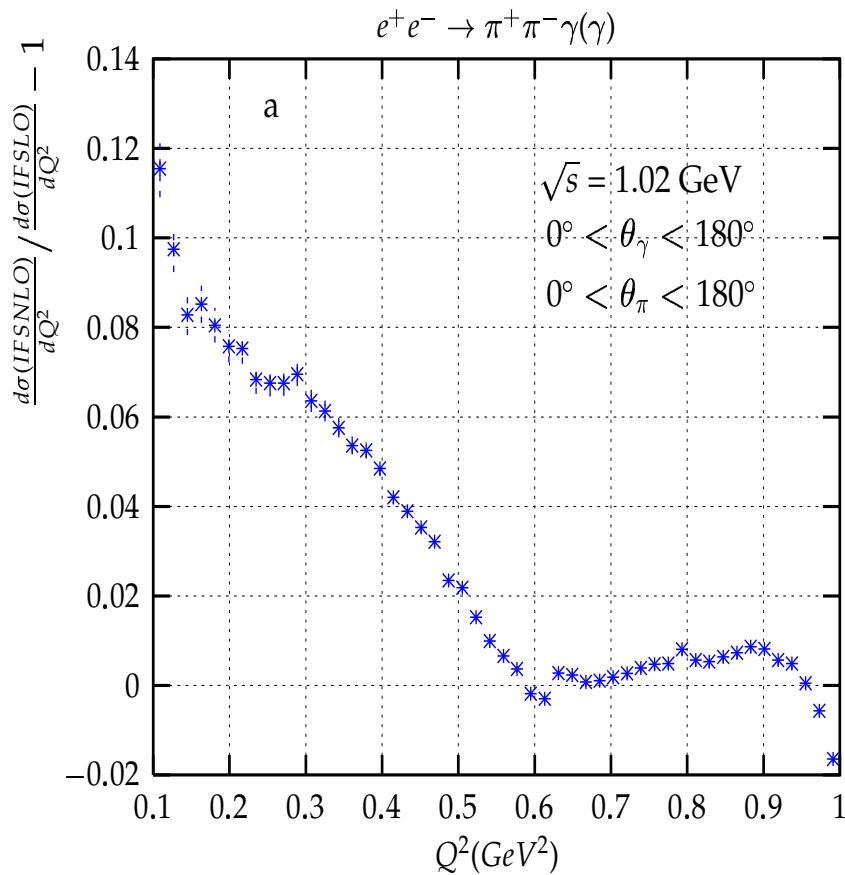


⇒ dominated by “two step process”:  $e^+e^- \rightarrow \gamma \rho (\rightarrow \gamma \pi\pi)$

⇒ importance of  $\pi\pi\gamma$  as input for  $a_\mu$



# Large effect for $Q^2 < m_\rho^2$ eliminated by suitable cuts on $\pi^+\pi^-$ configuration (suppress $2\gamma$ events)



⇒ Talk by D. Leone

or measure photon

# Experimental Perspectives

BABAR, BELLE

higher  $Q^2$  available

⇒ measurement of  $R(Q^2)$  from threshold up to at least 5 GeV.

Examples:

- ▶  $\pi\pi$
- ▶  $4\pi^\pm$
- ▶  $K K \pi\pi$
- ▶  $K K K K$

# PHOKHARA 4.0

- $\mu^+ \mu^- \gamma$  with FSR at NLO
- vacuum polarisation can be switched on
- nucleon pair production included

# To be done:

- three mesons:  $3\pi$  ( $\rightarrow \rho\pi$ ),  $KK\pi$
- $KK\pi\pi$ ,  $4K$
- narrow resonances

parameters of  $J/\psi$ ,  $\psi'$ :

observable:  $\Gamma_e \frac{\Gamma_f}{\Gamma_{tot}}$ ;  $f = \mu^+\mu^-, \pi^+\pi^-, 3\pi, 4\pi, 4K, \dots$

compare:  $\frac{\sigma_f}{\sigma_{\mu^+\mu^-}}(\text{off resonance}) \stackrel{?}{=} \frac{\sigma_f}{\sigma_{\mu^+\mu^-}}(\text{on resonance})$

$f = \mu^+\mu^-, \pi^+\pi^-, 4\pi, \dots$  virtual photon only ( $l=1$ )

$f = 3\pi, K\bar{K}, K\bar{K}\pi, \dots$  3 gluon intermediate state ( $l=0$ )

# III NUCLEON FORM FACTORS

(with Czyż, Nowak, Rodrigo, hep-ph/0403062)

$Q^2 \gtrsim 4m_N^2$  accessible at B-factories  
 $\Rightarrow$  study  $e^+e^- \rightarrow \gamma N\bar{N}$  (with  $N = p$  or  $n$ )

hadronic current:

$$J_\mu = -ie \cdot \bar{u}(q_2) \left( F_1^N(Q^2) \gamma_\mu - \frac{F_2^N(Q^2)}{4m_N} [\gamma_\mu, \not{Q}] \right) v(q_1),$$

$$Q = q_1 + q_2, \quad q = (q_1 - q_2)/2$$

or

$$G_M = F_1 + F_2, \quad G_E = F_1 + \frac{Q^2}{4m^2} F_2$$

## Result:

$$d\sigma = \frac{1}{2s} L_{\mu\nu} H^{\mu\nu} d\Phi_2(p_1 + p_2; Q, k) d\Phi_2(Q; q_1, q_2) \frac{dQ^2}{2\pi},$$

$$L_{\mu\nu} H^{\mu\nu} = \frac{(4\pi\alpha)^3}{Q^2} \left\{ \left( |G_M^N|^2 - \frac{1}{\tau} |G_E^N|^2 \right) \right. \\ \times \frac{32s}{\beta_N^2 (s - Q^2)} \left( \frac{1}{y_1} + \frac{1}{y_2} \right) \left( \frac{(p_1 \cdot q)^2 + (p_2 \cdot q)^2}{s^2} \right) \\ \left. + 2 \left( |G_M^N|^2 + \frac{1}{\tau} |G_E^N|^2 \right) \left[ \left( \frac{1}{y_1} + \frac{1}{y_2} \right) \frac{(s^2 + Q^4)}{s(s - Q^2)} - 2 \right] \right\},$$

where

$$y_{1,2} = \frac{s - Q^2}{2s} (1 \mp \cos \theta_\gamma), \quad \tau = \frac{Q^2}{4m_N^2}, \quad \beta_N^2 = 1 - \frac{4m_N^2}{Q^2}$$

## Separation of $|G_M|^2$ and $|G_E|^2$ through angular distribution:

$$L_{\mu\nu}H^{\mu\nu} = \frac{(4\pi\alpha)^3 (1 + \cos^2 \theta_\gamma)}{Q^2 (1 - \cos^2 \theta_\gamma)} \times 4 \left( |G_M^N|^2 (1 + \cos^2 \hat{\theta}) + \frac{1}{\tau} |G_E^N|^2 \sin^2 \hat{\theta} \right)$$

$\hat{\theta}$  = angle of nucleon with respect to  $\gamma$ -direction in hadronic rest frame

(valid for  $s/Q^2 \ll 1$ , corrections and “optimal frame”  $\rightarrow$  [hep-ph/0403062](https://arxiv.org/abs/hep-ph/0403062))

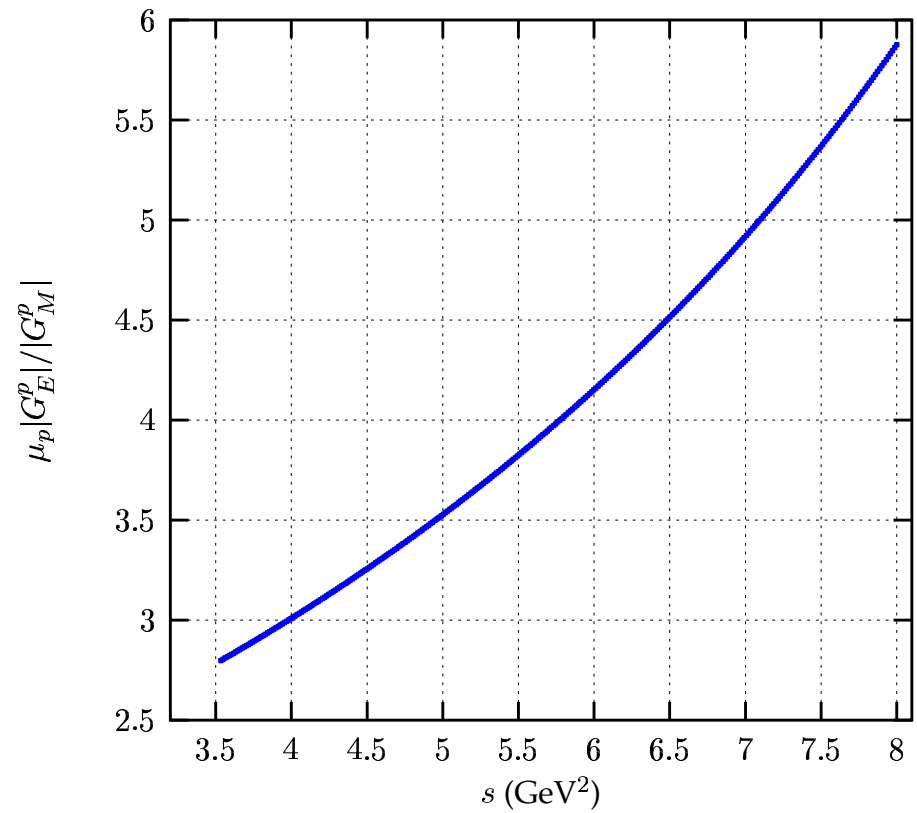
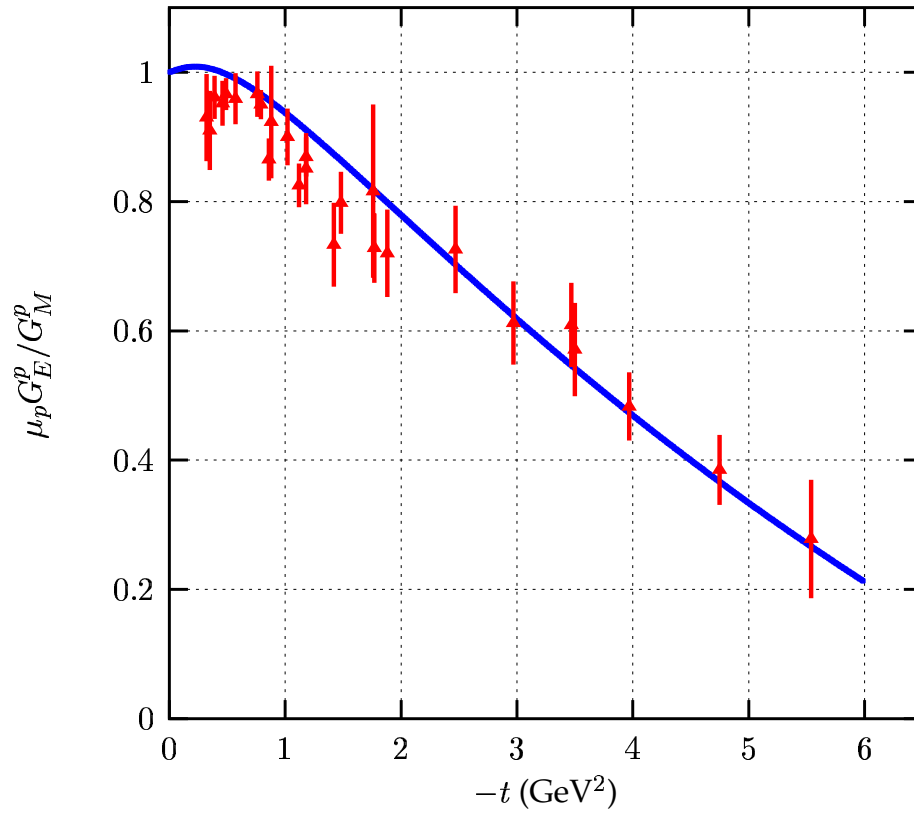
$\Rightarrow$  additional rotation by  $\theta_D = \frac{1}{2} \arctan \left( \frac{2s_\gamma c_\gamma}{\gamma (\beta^2 + c_\gamma^2 - s_\gamma^2/\gamma^2)} \right) \approx \frac{1}{\gamma} \frac{s_\gamma c_\gamma}{1 + c_\gamma^2}$

with  $s_\gamma = \sin \theta_\gamma$ ,  $\beta = (s - Q^2)/(s + Q^2)$ ,  $\gamma = (s + Q^2)/2\sqrt{sQ^2}$

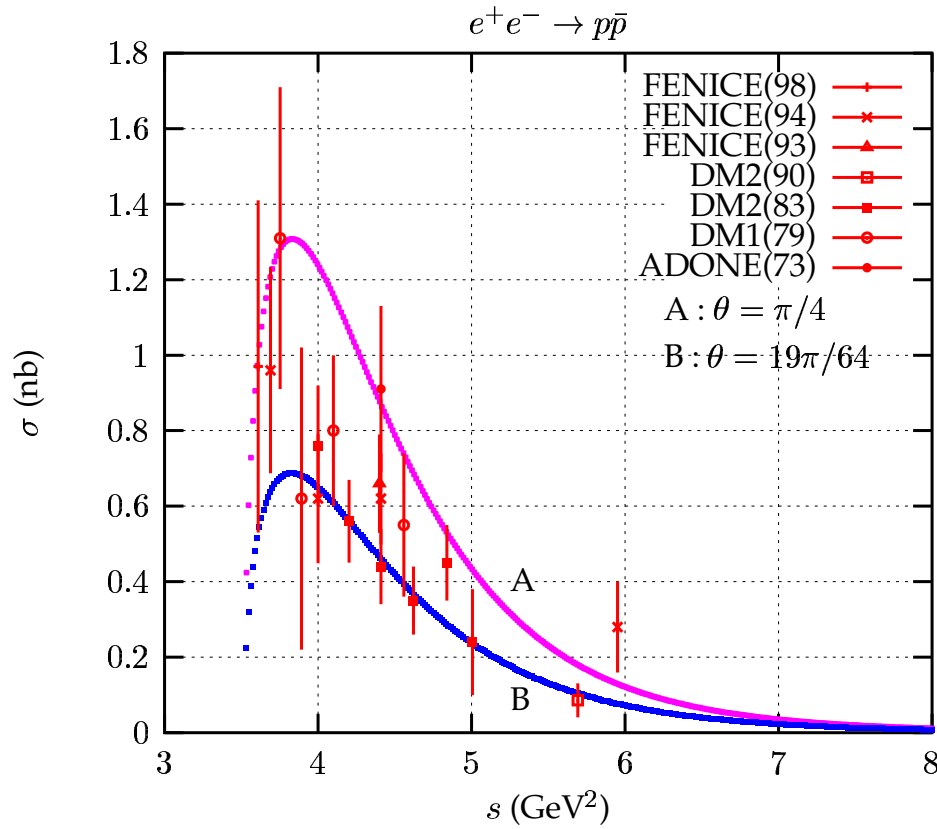
**Similarity to  $e^+e^- \rightarrow N\bar{N}$ :**

$$\frac{d\sigma}{d\Omega} = \frac{\alpha^2 \beta_N}{4Q^2} \left( |G_M^N|^2 (1 + \cos^2 \theta) + \frac{1}{\tau} |G_E^N|^2 \sin^2 \theta \right)$$

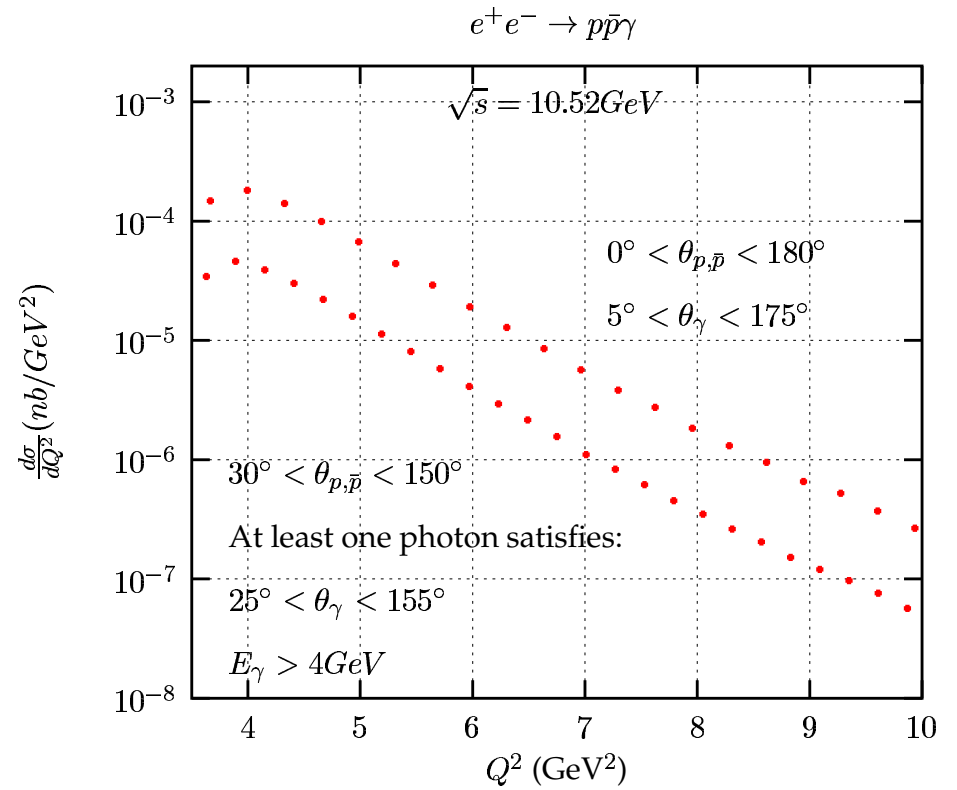
# Implementation on basis of model for form factor:







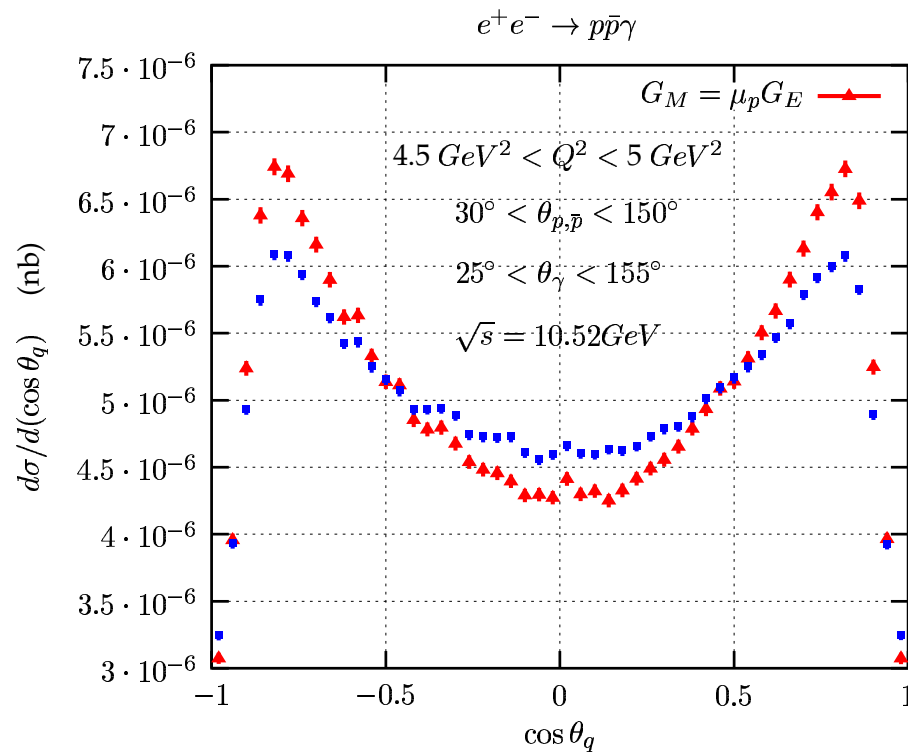
$$e^+e^- \rightarrow p\bar{p}$$



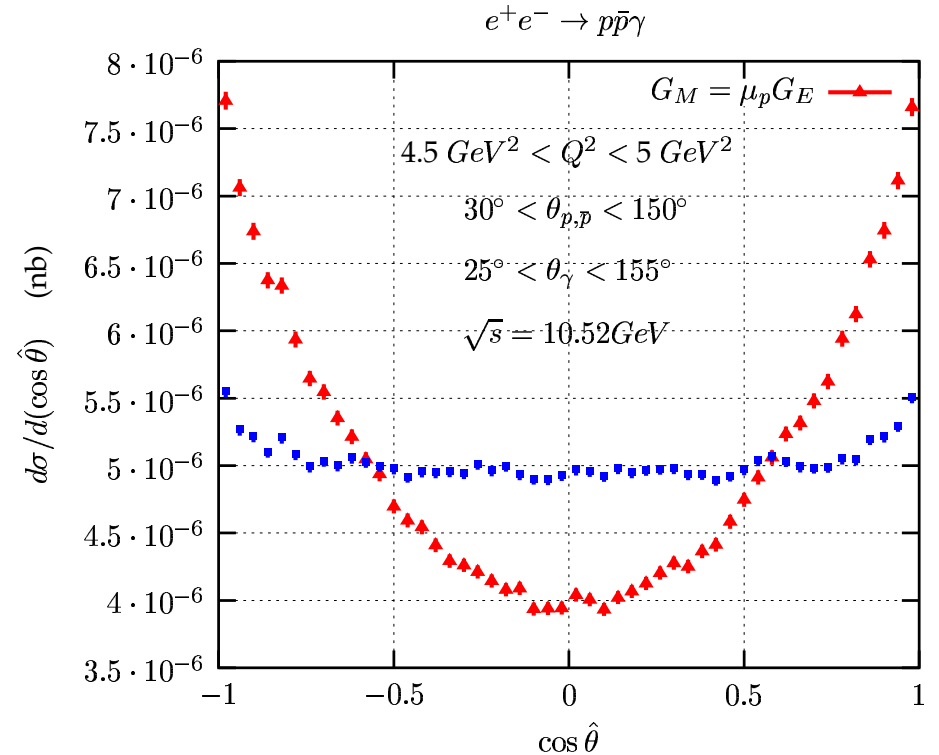
$$e^+e^- \rightarrow p\bar{p}\gamma$$

implementation in PHOKHARA

# Angular distributions of nucleon



lab frame



hadronic rest frame

(two choices for  $G_M/G_E$ )

# Comments

- similar results for **neutron** pair production
- NLO corrections from ISR included (corrections  $\sim 1\text{--}2\%$ )
- no FSR

thousands of events around  $4\text{--}5 \text{ GeV}^2$   
several events up to  $7\text{--}8 \text{ GeV}^2$

# IV MESON FORM FACTORS at LARGE $Q^2$

(with Bruch, Khodjamirian, hep-ph/0409080)

radiative return will explore large  $Q^2$

convenient representation for  $F_\pi$  :

generalized VDM with  $\rho, \rho', \dots$

combined with Veneziano-type tower of resonances (Dominguez)

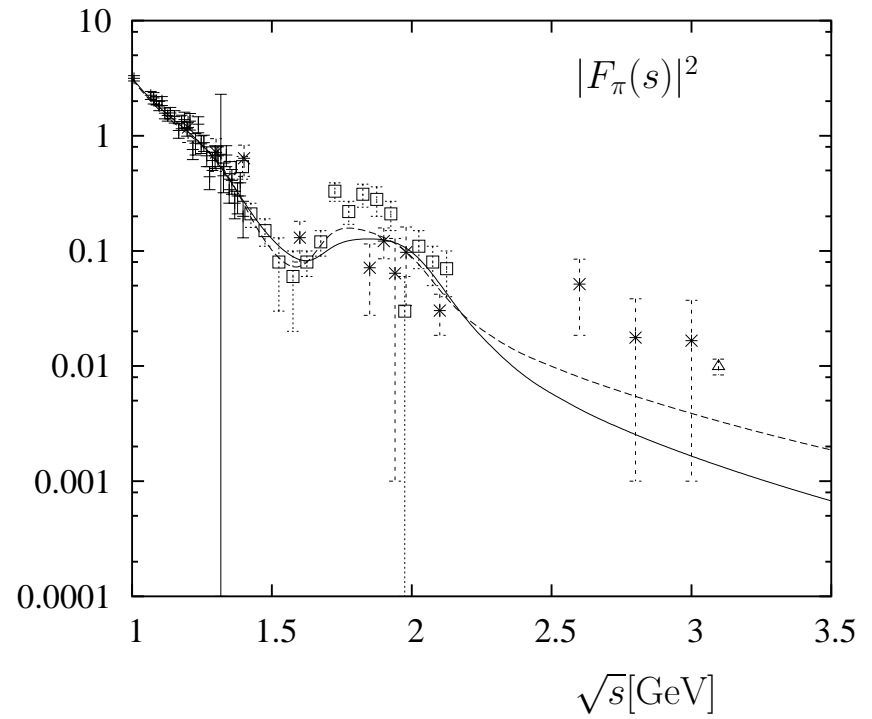
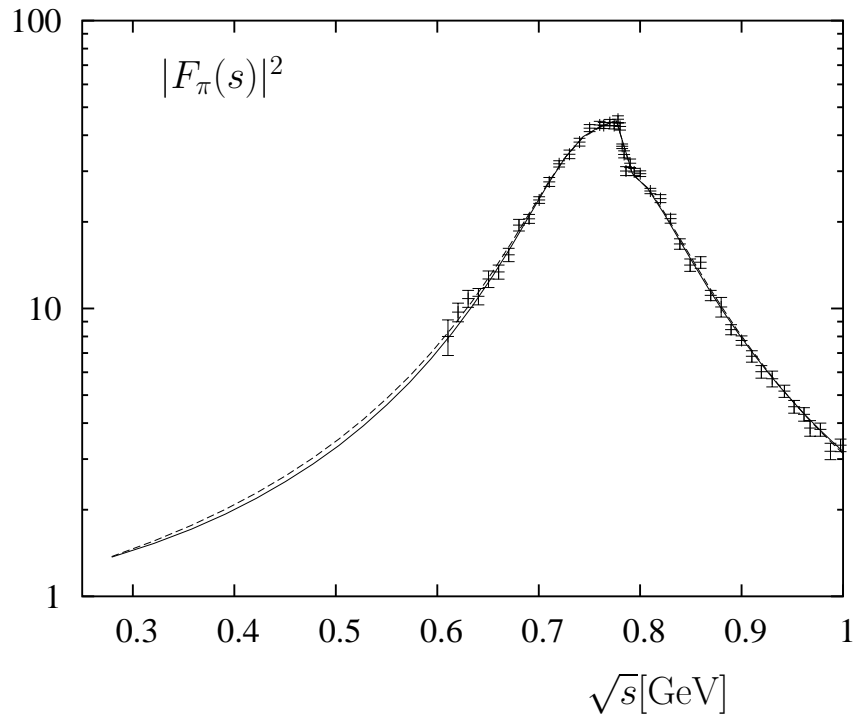
$$F_\pi(s) = \sum_{n=0}^{\infty} c_n \frac{m_n^2}{m_n^2 - s},$$
$$c_n = \frac{(-1)^n \Gamma(\beta - 1/2)}{\sqrt{\pi} (\frac{1}{2} + n) \Gamma(n + 1) \Gamma(\beta - 1 - n)},$$
$$m_n^2 = m_\rho^2 (1 + 2n),$$

$\beta$  = free parameter

# Modifications:

- finite widths
- parameters of  $\rho$ ,  $\rho'$ ,  $\rho''$  fitted to data
- Breit-Wigner for  $\rho$ ,  $\rho'$ ,  $\rho''$  with  $Q^2$ -dependent widths  
 $\Rightarrow$  reasonable agreement between model and fit

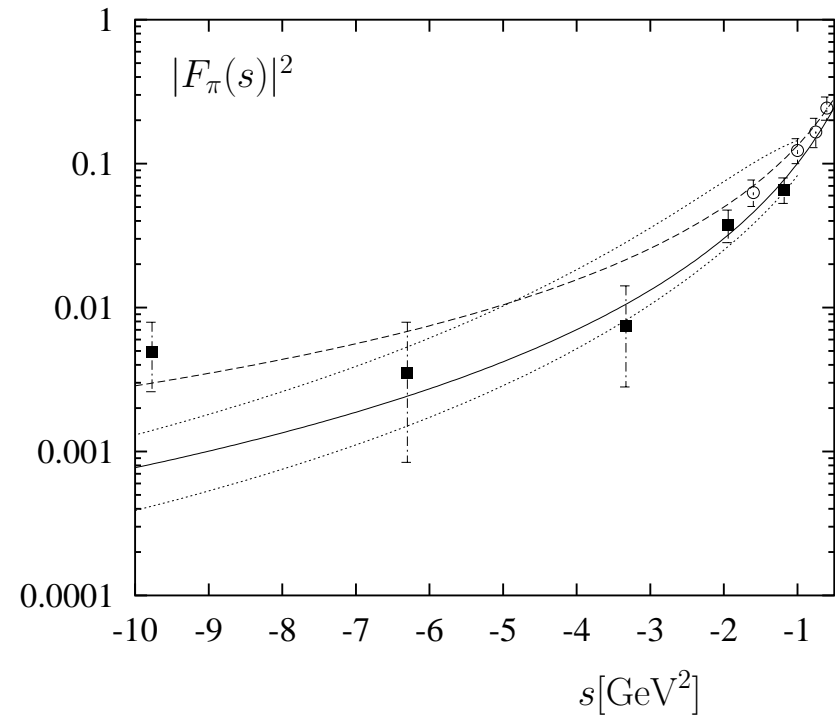
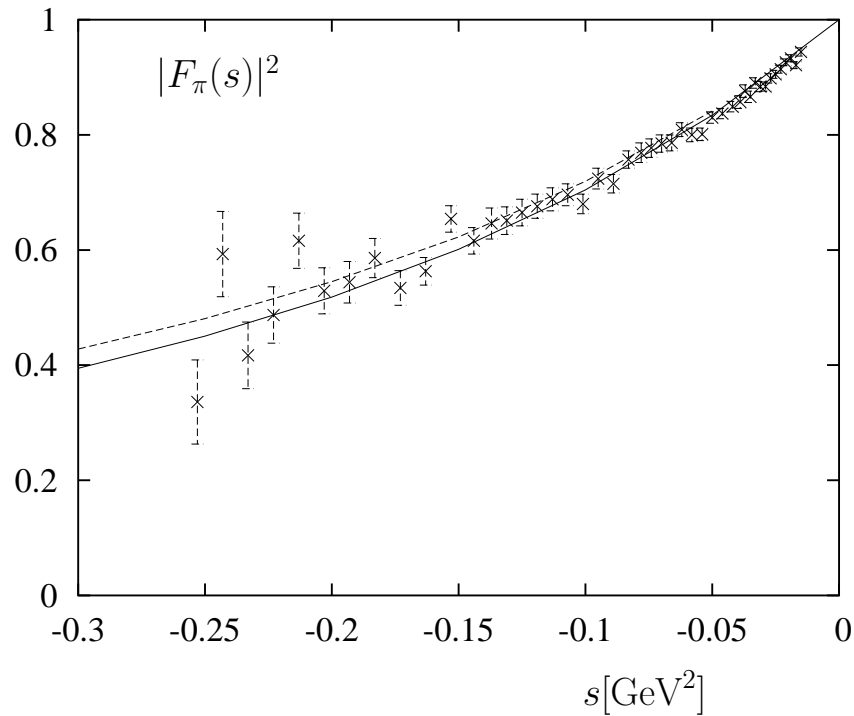
Parameter	Input	Fit(KS)	Fit(GS)	dual- QCD $_{N_c=\infty}$	PDG value
$m_\rho$	-	$773.9 \pm 0.6$	$776.3 \pm 0.6$	input	$775.5 \pm 0.5$
$\Gamma_\rho$	-	$144.9 \pm 1.0$	$150.5 \pm 1.0$	input	$150.3 \pm 1.6$
$m_\omega$	783.0	-	-	-	$782.59 \pm 0.11$
$\Gamma_\omega$	8.4	-	-	-	$8.49 \pm 0.08$
$m_{\rho'}$	-	$1357 \pm 18$	$1380 \pm 18$	1335	$1465 \pm 25$
$\Gamma_{\rho'}$	-	$437 \pm 60$	$340 \pm 53$	266	$400 \pm 60$
$m_{\rho''}$	1700	-	-	1724	$1720 \pm 20$
$\Gamma_{\rho''}$	240	-	-	344	$250 \pm 100$
$m_{\rho'''}$	-	-	-	2040	-
$\Gamma_{\rho'''}$	-	-	-	400	-
$c_0$	-	$1.171 \pm 0.007$	$1.098 \pm 0.005$	1.171	-
$\beta$	$c_0$	$2.30 \pm 0.01$	$2.16 \pm 0.015$	2.3(input)	-
$c_\omega$	0.00184(KS) 0.00195(GS)	-	-	-	-
$c_1$	-	$-0.119 \pm 0.011$	$-0.069 \pm 0.009$	-0.1171	-
$c_2$	-	$0.0115 \pm 0.0064$	$0.0216 \pm 0.0064$	-0.0246	-
$c_3$	$\sum c_n=1$	$-0.0438 \mp 0.02$	$-0.0309 \mp 0.02$	-0.00995	-
$\sum_{n=4}^{\infty} c_n$	-0.01936	-	-	-0.01936	-
$\chi^2/d.o.f.$	-	155/101	153/101	-	-



data point at 3.1 GeV ( $J/\Psi \rightarrow \pi\pi$ ) cannot be accommodated

# spacelike region:

good agreement with data and with sum rules



$$e^+e^- \rightarrow K^+K^-, K^0\bar{K}^0$$

isospin symmetry:

$$F_{K^+} = +F^{(I=1)} + F^{(I=0)}$$

$$F_{K^0} = -F^{(I=1)} + F^{(I=0)}$$

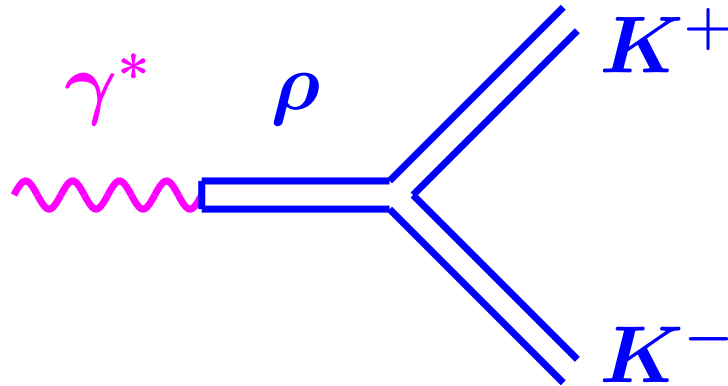
resonances:

$$\begin{aligned} F_{K^+}(s) = & +\frac{1}{2} \left( c_{\rho}^K BW_{\rho}(s) + c_{\rho'}^K BW_{\rho'}(s) + c_{\rho''}^K BW_{\rho''}(s) \right) \\ & + \frac{1}{6} \left( c_{\omega}^K BW_{\omega}(s) + c_{\omega'}^K BW_{\omega'}(s) + c_{\omega''}^K BW_{\omega''}(s) \right) \\ & + \frac{1}{3} \left( c_{\phi} BW_{\phi}(s) + c_{\phi'} BW_{\phi'}(s) \right), \end{aligned}$$

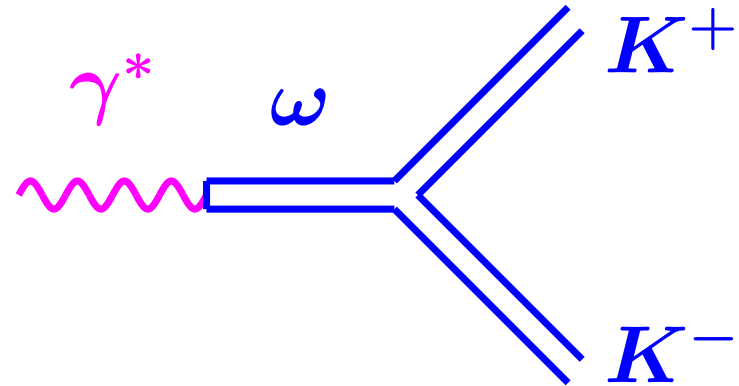
$$\begin{aligned} F_{K^0}(s) = & -\frac{1}{2} \left( c_{\rho}^K BW_{\rho}(s) + c_{\rho'}^K BW_{\rho'}(s) + c_{\rho''}^K BW_{\rho''}(s) \right) \\ & + \frac{1}{6} \left( c_{\omega}^K BW_{\omega}(s) + c_{\omega'}^K BW_{\omega'}(s) + c_{\omega''}^K BW_{\omega''}(s) \right) \\ & + \frac{1}{3} \left( \eta_{\phi} c_{\phi} BW_{\phi}(s) + c_{\phi'} BW_{\phi'}(s) \right) \end{aligned}$$



# quark model:



$$\frac{1}{\sqrt{2}} f_\rho \quad g_{\rho KK}$$



$$\frac{1}{3\sqrt{2}} f_\omega \quad g_{\omega KK}$$

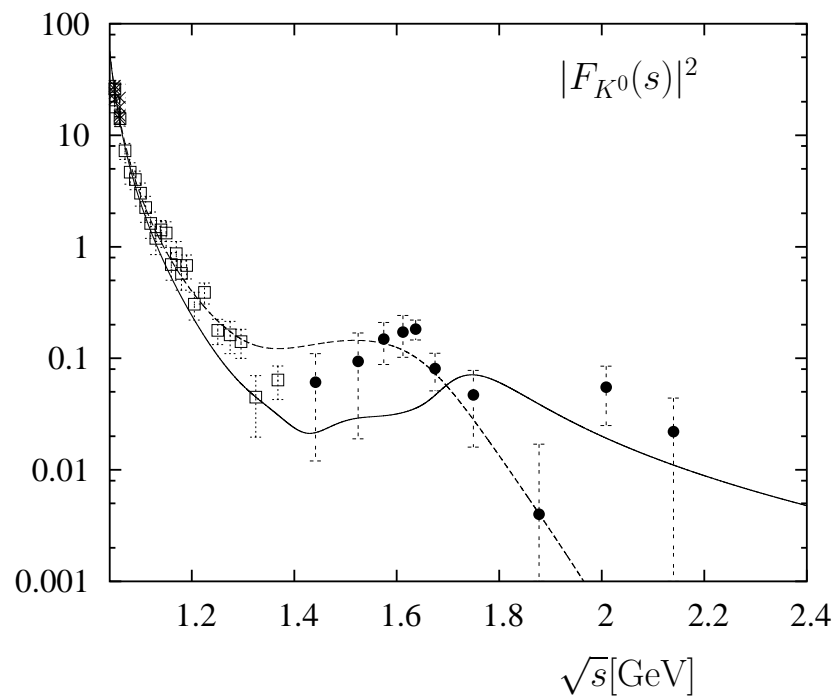
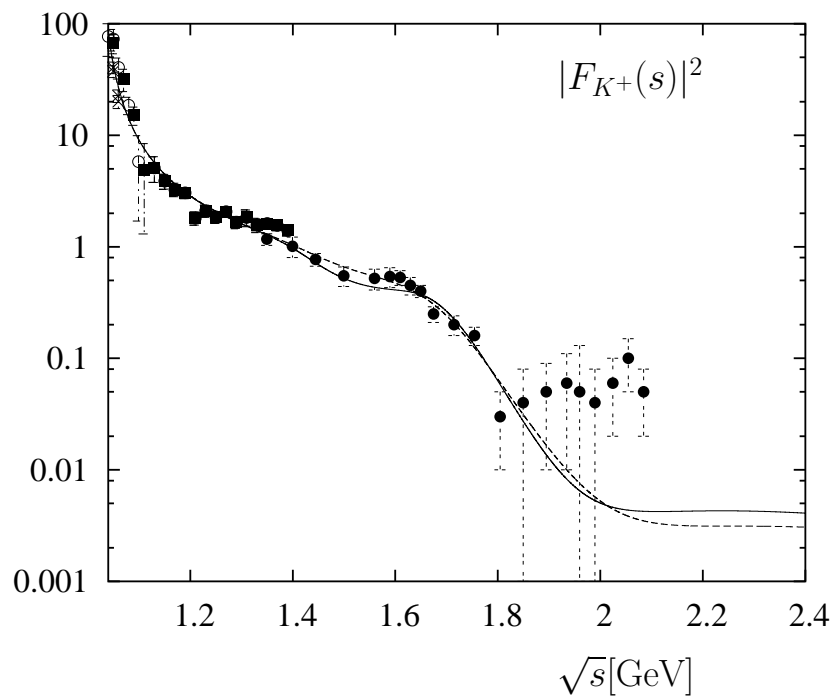
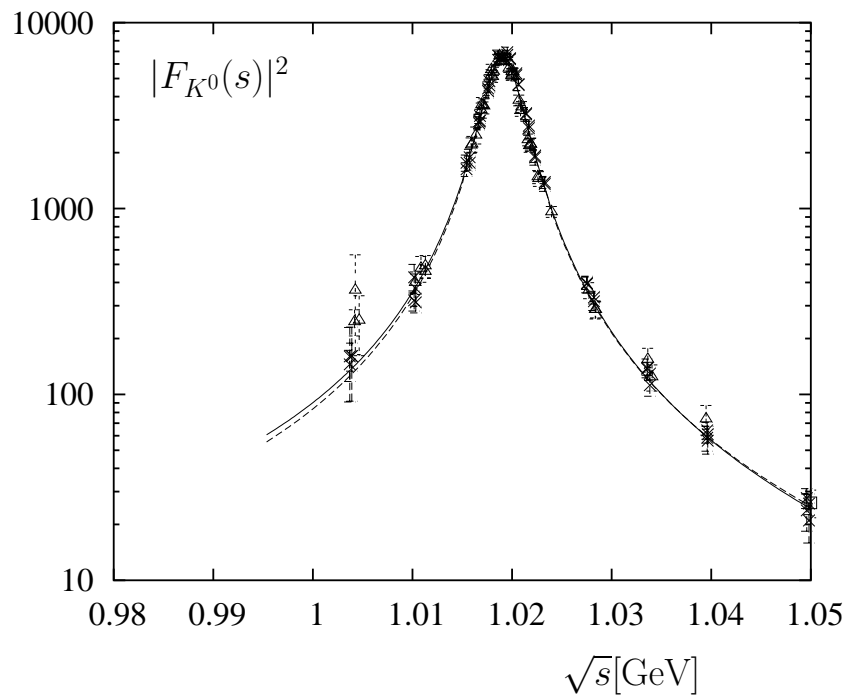
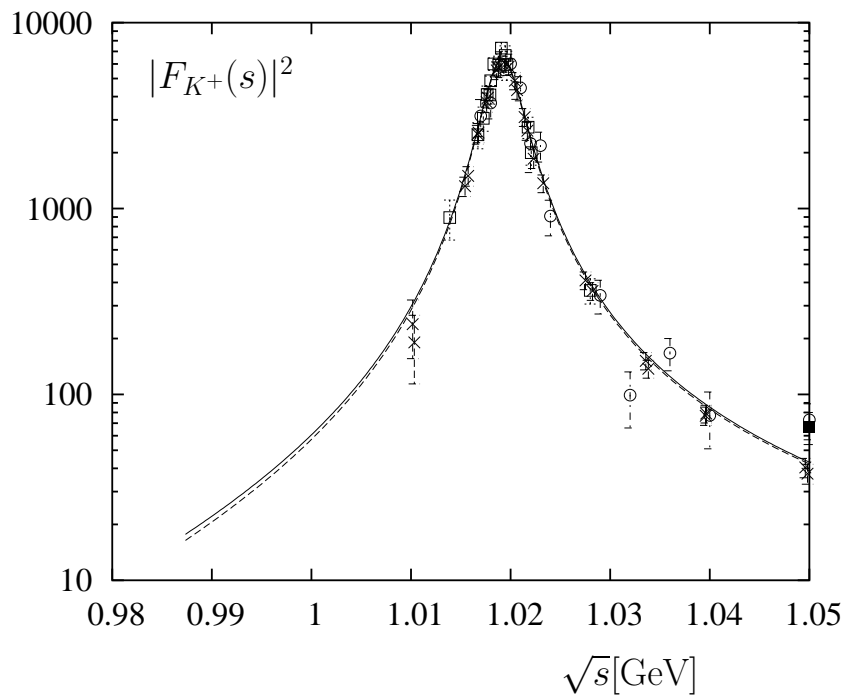
$$\text{constraint: } f_\rho = f_\omega, \quad g_{\rho KK} = g_{\omega KK}$$

$$\Rightarrow c_\rho = c_\omega$$

fit performed with **(solid curves)**  
or **without (dashed curves)** this constraint

# Results:

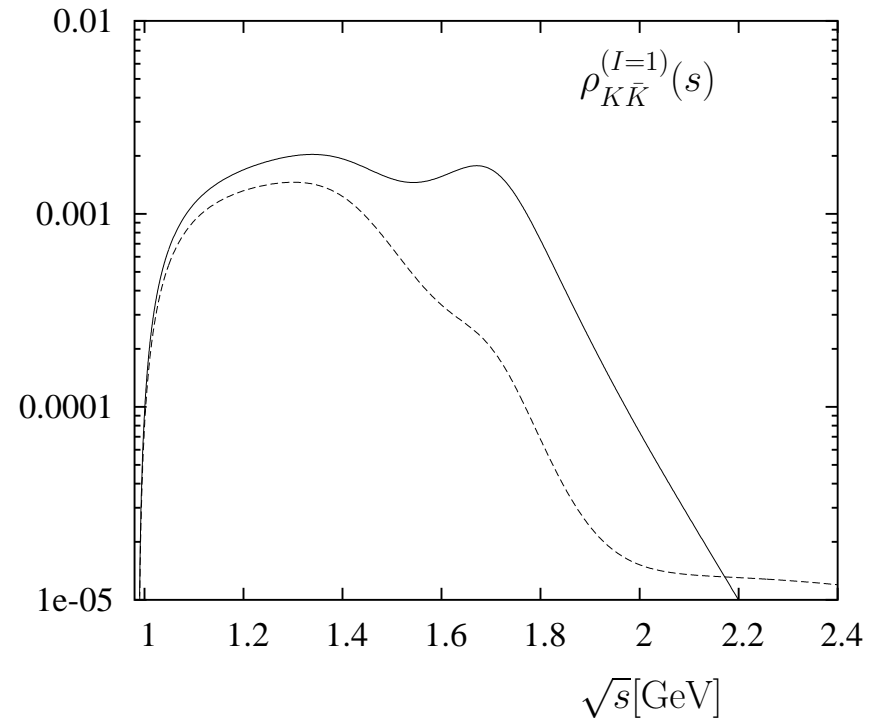
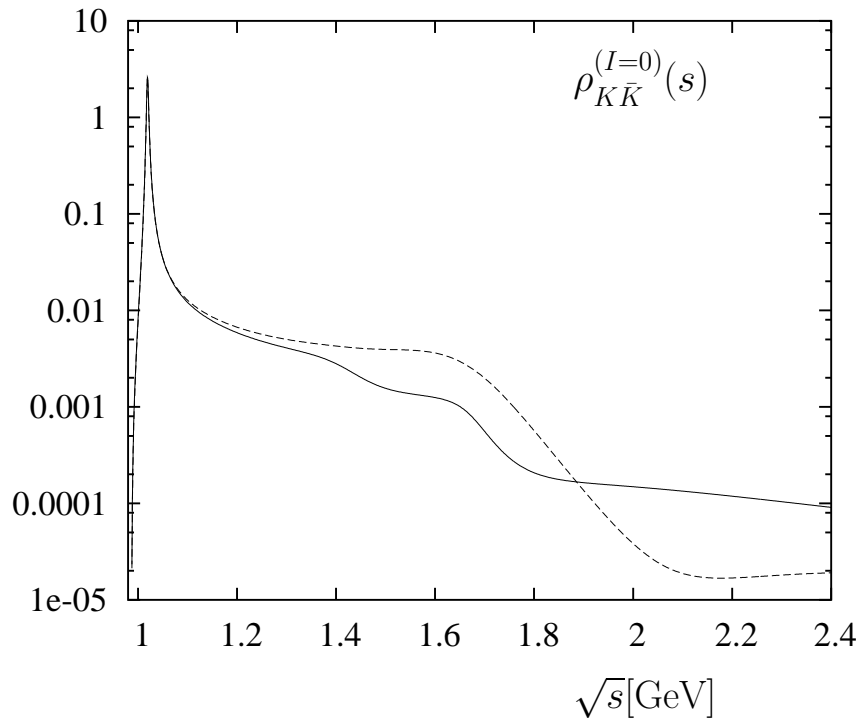
Parameter	Input	Fit(1)	Fit(2)	PDG value
$m_\phi$	-	$1019.372 \pm 0.02$	$1019.355 \pm 0.02$	$1019.456 \pm 0.02$
$\Gamma_\phi$	-	$4.36 \pm 0.05$	$4.29 \pm 0.05$	$4.26 \pm 0.05$
$m_{\phi'}$	1680	-	-	$1680 \pm 20$
$\Gamma_{\phi'}$	150	-	-	$150 \pm 50$
$m_\rho$	775	-	-	$775.8 \pm 0.5$
$\Gamma_\rho$	150	-	-	$150.3 \pm 1.6$
$m_{\rho'}$	1465	-	-	$1465 \pm 25$
$\Gamma_{\rho'}$	400	-	-	$400 \pm 60$
$m_{\rho''}$	1720	-	-	$1720 \pm 20$
$\Gamma_{\rho''}$	250	-	-	$250 \pm 100$
$m_\omega$	783.0	-	-	$782.59 \pm 0.11$
$\Gamma_\omega$	8.4	-	-	$8.49 \pm 0.08$
$m_{\omega'}$	1425	-	-	1400-1450
$\Gamma_{\omega'}$	215	-	-	180-250
$m_{\omega''}$	1670	-	-	$1670 \pm 30$
$\Gamma_{\omega''}$	315	-	-	$315 \pm 35$
$c_\phi$	-	$1.018 \pm 0.006$	$0.999 \pm 0.007$	-
$c_{\phi'}$	$1 - c_\phi^K$	$-0.018 \mp 0.006$	$0.001 \mp 0.007$	-
$c_\rho^K$	-	$1.195 \pm 0.009$	$1.139 \pm 0.010$	-
$c_{\rho'}^K$	-	$-0.112 \pm 0.010$	$-0.124 \pm 0.012$	-
$c_{\rho''}^K$	$1 - c_\rho^K - c_{\rho'}^K$	$-0.083 \mp 0.019$	$-0.015 \mp 0.022$	-
$c_\omega^K(1)$	$c_\rho^K$	$1.195 \pm 0.009$	-	-
$c_\omega^K(2)$	-	-	$1.467 \pm 0.035$	-
$c_{\omega'}^K(1)$	$c_{\rho'}^K$	$-0.112 \pm 0.010$	-	-
$c_{\omega'}^K(2)$	-	-	$-0.018 \pm 0.024$	-
$c_{\omega''}^K$	$1 - c_\omega^K - c_{\omega'}^K$	$-0.083 \mp 0.019$	$-0.449 \mp 0.059$	-
$\chi^2/d.o.f.$	-	328/242	281/240	-



## Spectral function separated for $I = 0$ and $I = 1$

(useful for electroweak analysis!)

$$\rho_{K\bar{K}}^{(I=0,1)}(s) = \frac{1}{12\pi} \left| \frac{F_{K^+}(s) \pm F_{K^0}(s)}{2} \right|^2 \left( \frac{2p_K(s)}{\sqrt{s}} \right)^3$$



significant model dependence above 1.5 GeV (poor data for  $|F_{K^0}|^2$  !)

$$\tau \rightarrow K^- K^0 \nu$$

Predictions based on isospin symmetry and  $I = 1$  part of form factor:

$$\left( \frac{1}{BR(\tau \rightarrow \mu^- \bar{\nu}_\mu \nu_\tau)} \right) \frac{dBR(\tau \rightarrow K^- K^0 \nu_\tau)}{d\sqrt{Q^2}} =$$

$$\frac{|V_{ud}|^2}{2m_\tau^2} \left( 1 + \frac{2Q^2}{m_\tau^2} \right) \left( 1 - \frac{Q^2}{m_\tau^2} \right)^2 \left( 1 - \frac{4m_K^2}{Q^2} \right)^{3/2}$$

$$\times \sqrt{Q^2} |F_{K^- K^0}(Q^2)|^2$$

and  $F_{K^- K^0} = -F_{K^+} + F_{K^0}$

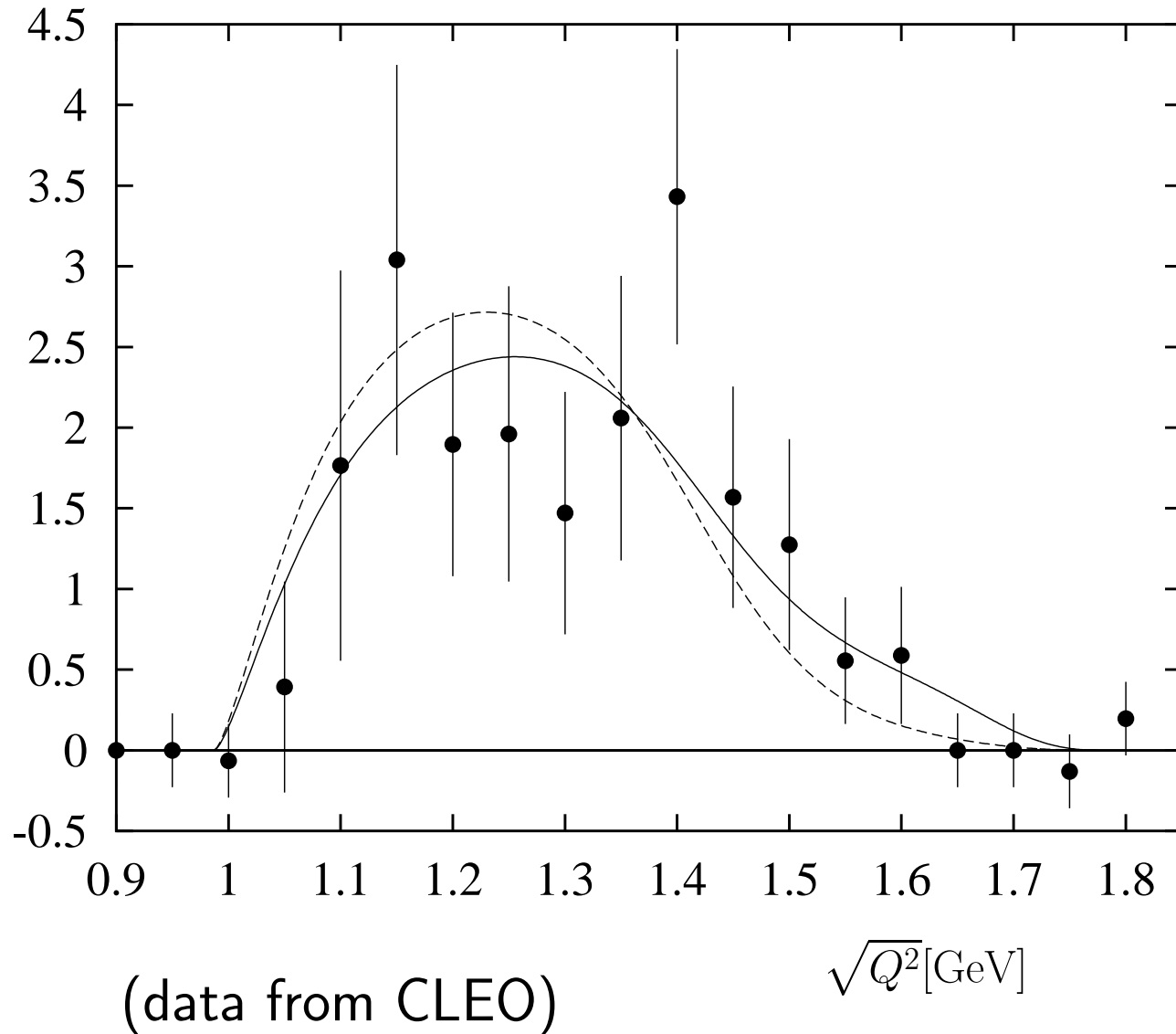
$$\Rightarrow BR(\tau \rightarrow K^- K^0 \nu_\tau) = 0.19 \pm 0.01\% \quad (0.13 \pm 0.01\%)$$

to be compared with

$$BR(\tau \rightarrow K^- K^0 \nu_\tau) = 0.154 \pm 0.016\%.$$

# $Q^2$ distribution:

will provide further constraints!



# V Conclusions

- continuous development of **PHOKHARA**
  - ⇒ radiative corrections
  - ⇒ more channels
  - ⇒ cooperation between theory and experiment crucial
- nucleon form factors:  
 $G_E$  and  $G_M$  can be measured for a wide range of  $Q^2$
- pion form factor: structures at large  $Q^2$   
kaon form factors:  $K^+K^-$  &  $K^0\bar{K}^0 \Rightarrow K^-K^0$   
⇒ prediction for  $\tau \rightarrow \nu K^- K^0$