

## B PHYSICS AND CP VIOLATION

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- **U-spin: Equal CP rate asymmetries in pairs of charmless decays** (M.G.)
- **Determination of weak phases**
  - $\phi_3 \equiv \gamma$  from U-spin  
 $B^0 \rightarrow K^\pm \pi^\mp, B_s \rightarrow K^\mp \pi^\pm$  (M.G. & J. Rosner)
  - Flavor SU(3): What's special about  
$$\frac{2\Gamma(B^\pm \rightarrow K^\pm \pi^0)}{\Gamma(B^\pm \rightarrow K^0 \pi^\pm)} = 1.45 \pm 0.46 \neq 1 ?$$
(M. Neubert & J. Rosner)
  - $\phi_2 \equiv \alpha$  from  $B \rightarrow \pi\pi$   
isospin analysis and alternatives (D. London)
  - Other determinations of  $\phi_3 \equiv \gamma$   
 $B^\pm \rightarrow DK^\pm$  (D. Atwood, S. Oh)  
 $B_s(t) \rightarrow D_s K$  (A. Falk)

## Theorem about equal CP rate differences

$$\mathcal{H}_{\text{eff}}^{\Delta S=1} \equiv \frac{G_F}{\sqrt{2}} V_{ub}^* V_{us} \left( \sum_1^2 c_i Q_i^{us} + \sum_3^{10} c_i Q_i^s \right) \\ + V_{cb}^* V_{cs} \left( \sum_1^2 c_i Q_i^{cs} + \sum_3^{10} c_i Q_i^s \right)$$

U-Spin Doublets  $\begin{pmatrix} d \\ s \end{pmatrix}$   $\begin{matrix} \bar{b}u\bar{u}s \\ \bar{b}c\bar{c}s \end{matrix}$   $\bar{b}_s \sum_{q'} (e_{q'}) \bar{q}' q'$

$$\mathcal{H}_{\text{eff}}^{\Delta S=1} = V_{ub}^* V_{us} U^s + V_{cb}^* V_{cs} C^s$$

$$s \rightarrow d: \quad \mathcal{H}_{\text{eff}}^{\Delta S=0} = V_{ub}^* V_{ud} U^d + V_{cb}^* V_{cd} C^d$$

Decays related by  $U(d \leftrightarrow s)$

$$\langle f | U^s (C^s) | B \rangle \equiv \langle Uf | U^d (C^d) | UB \rangle \equiv A_u (A_c)$$

$$A(B \rightarrow f, \Delta S = 1) = V_{ub}^* V_{us} A_u + V_{cb}^* V_{cs} A_c$$

$$A(UB \rightarrow Uf, \Delta S = 0) = V_{ub}^* V_{ud} A_u + V_{cb}^* V_{cd} A_c$$

$$A(\bar{B} \rightarrow \bar{f}, \Delta S = -1) = V_{ub} V_{us}^* A_u + V_{cb} V_{cs}^* A_c$$

$$A(U\bar{B} \rightarrow U\bar{f}, \Delta S = 0) = V_{ub} V_{ud}^* A_u + V_{cb} V_{cd}^* A_c$$

$$\frac{V_{ub}^* V_{us}}{V_{ub}^* V_{ud}} \equiv -\frac{V_{cb}^* V_{cd}}{V_{cb}^* V_{cs}} \equiv \tan \theta_c$$

Question : 4 rates determine  $|V_{ub}^* V_{us} A_u|$ ,  $|V_{cb}^* V_{cs} A_c|$   
 $\delta \equiv \text{Arg}(A_u A_c^*)$ ,  $\phi_3 \equiv \gamma \equiv \text{Arg}(-V_{ub}^* V_{ud} V_{cb} V_{cd}^*)$  ???

Answer: Needs additional input

CKM unitarity

$$\text{Im}(V_{ub}^* V_{us} V_{cb} V_{cs}^*) \equiv -\text{Im}(V_{ub}^* V_{ud} V_{cb} V_{cd}^*)$$

implies

**Theorem:** in the U-spin symmetry limit

$$\begin{aligned} \Delta\Gamma(B \rightarrow f) &\equiv |A(B \rightarrow f)|^2 - |A(\bar{B} \rightarrow \bar{f})|^2 \equiv \\ -\Delta\Gamma(UB \rightarrow Uf) &\equiv -[|A(UB \rightarrow Uf)|^2 - |A(U\bar{B} \rightarrow U\bar{f})|^2] \end{aligned}$$

examples:

$$\Delta\Gamma(B^0 \rightarrow K^+ \pi^-) \equiv -\Delta\Gamma(B_s \rightarrow \pi^+ K^-)$$

$$\Delta\Gamma(B^0 \rightarrow K^{*+} \pi^-) \equiv -\Delta\Gamma(B_s \rightarrow \rho^+ K^-)$$

$$\Delta\Gamma(B^+ \rightarrow K^+ \pi^+ \pi^-) \equiv -\Delta\Gamma(B^+ \rightarrow \pi^+ K^+ K^-)$$

$$\Delta\Gamma(B^+ \rightarrow K^{*+} \gamma) \equiv -\Delta\Gamma(B^+ \rightarrow \rho^+ \gamma)$$

CKM  $\Rightarrow$  opposite sign CP asymmetries

**U-spin: determining  $\phi_3 \equiv \gamma$**

$B^0 \rightarrow K^+ \pi^-, B_s \rightarrow K^- \pi^+, B^+ \rightarrow K^0 \pi^+, \text{ c.c.}$

$$A(B^0 \rightarrow K^+ \pi^-) \equiv V_{ub}^* V_{us} A_u + V_{cb}^* V_{cs} A_c$$

$$A(B_s \rightarrow \pi^+ K^-) \equiv V_{ub}^* V_{ud} A_u + V_{cb}^* V_{cd} A_c$$

$$A(\bar{B}^0 \rightarrow K^- \pi^+) \equiv V_{ub} V_{us}^* A_u + V_{cb} V_{cs}^* A_c$$

$$A(\bar{B}_s \rightarrow \pi^- K^+) \equiv V_{ub} V_{ud}^* A_u + V_{cb} V_{cd}^* A_c$$

$$\begin{aligned} \Gamma(B_s \rightarrow K^- \pi^+) - \Gamma(\bar{B}_s \rightarrow K^+ \pi^-) &\equiv \\ -[\Gamma(B^0 \rightarrow K^+ \pi^-) - \Gamma(\bar{B}^0 \rightarrow K^- \pi^+)] &\equiv \end{aligned}$$

**additional input for  $\gamma$ :  $\Gamma(B^\pm \rightarrow K_S \pi^\pm)$**

neglect **rescattering** (no "Tree"  $\propto V_{ub}^* V_{us}$ )

$$A(B^+ \rightarrow K^0 \pi^+) \equiv A(B^- \rightarrow \bar{K}^0 \pi^-) \equiv V_{cb}^* V_{cs} A_c$$

$$R \equiv \frac{\Gamma(B^0 \rightarrow K^\pm \pi^\mp)}{\Gamma(B^\pm \rightarrow K \pi^\pm)} \quad R_s \equiv \frac{\Gamma(B_s \rightarrow K^\pm \pi^\mp)}{\Gamma(B^\pm \rightarrow K \pi^\pm)}$$

$$\mathcal{A}_0 \equiv \frac{\Delta(B^0 \rightarrow K^+ \pi^-)}{\Gamma(B^\pm \rightarrow K \pi^\pm)} \quad \mathcal{A}_s \equiv \frac{\Delta(B_s \rightarrow K^- \pi^+)}{\Gamma(B^\pm \rightarrow K \pi^\pm)}$$

$$r \equiv |V_{ub}^* V_{us} A_u| / |V_{cb}^* V_{cs} A_c|$$

$$\delta_{K\pi} \equiv \text{Arg}(A_u A_c^*)$$

$$\gamma \equiv \text{Arg}(-V_{ub}^* V_{ud} V_{cb} V_{cd}^*)$$

$$R \equiv 1 + r^2 + 2r \cos \delta_{K\pi} \cos \gamma$$

$$R_s \equiv \tan^2 \theta_c^2 + (r / \tan \theta_c)^2 - 2r \cos \delta_{K\pi} \cos \gamma$$

$$\mathcal{A}_0 \equiv -\mathcal{A}_s \equiv -2r \sin \delta_{K\pi} \sin \gamma \quad \text{checks } SU(3)$$

*SU(3) breaking (factorization):*

$$f \equiv F_{B_s K}(m_\pi^2) f_\pi / F_{B\pi}(m_K^2) f_K$$

$$\mathcal{A}_s \equiv -f^2 \mathcal{A}_0$$

Tevatron Run II:  $\delta(\gamma) \sim 10^\circ$

## Another example of U-spin for

$$\phi_3 \equiv \gamma \quad (R.Fleischer)$$

$$B^0(t) \rightarrow \pi^+ \pi^-, \quad B_s(t) \rightarrow K^+ K^-$$

measure

$$A_{sym}(t) \equiv A_{mix} \sin(\Delta mt) + A_{dir} \cos(\Delta mt)$$

for  $B^0 \rightarrow \pi^+ \pi^-$  and  $B_s \rightarrow K^+ K^-$

Four quantities in terms of  $|P/T|_{\pi\pi}$ ,  $\delta$ ,  $\beta$ ,  $\gamma$ .

$$\text{Tevatron Run II: } \delta(\gamma) \sim 10^\circ$$

## Flavor SU(3) in $B, B_s \rightarrow \pi\pi, K\pi, K\bar{K}$ (rescattering)

$$A \equiv \langle (PP)_{S\text{-wave}} | \text{effective Hamiltonian} | B_{u,d,s} \rangle$$

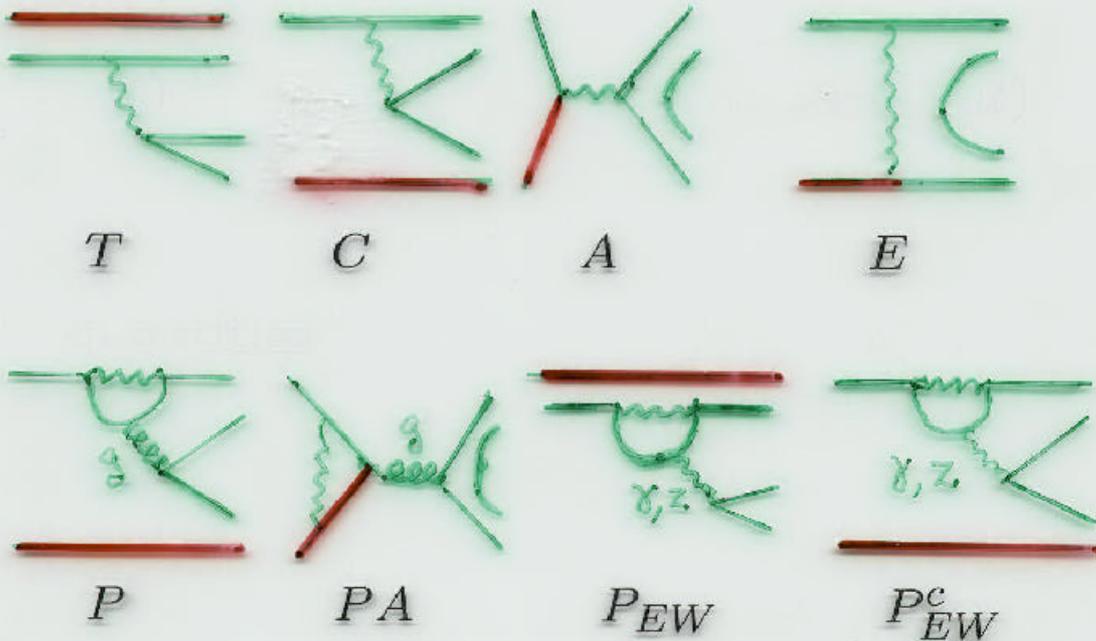
$$(8 \otimes 8)_{\text{sym}} \equiv 1 + 8 + 27 \quad \bar{3} + 6 + \bar{15} \quad 3$$

5 SU(3) reduced matrix elements:

$$\langle 1 || \bar{3} || 3 \rangle \quad \langle 8 || \bar{3} || 3 \rangle \quad \langle 8 || 6 || 3 \rangle \quad \langle 8 || \bar{15} || 3 \rangle \quad \langle 27 || \bar{15} || 3 \rangle$$

equivalent to 8 diagrams in 5 combinations

occurring with given CKM factors



## EW penguin $\propto$ Tree

- "Tree" operators,  $Q_1$  and  $Q_2$ :  $(V - A)(V - A)$
- EW penguin operators with **dominant**  $c_i$ ,  $Q_9$  and  $Q_{10}$ :  $(V - A)(V - A)$

$\Rightarrow$  EWP and Tree operators transforming as given SU(3) representations are proportional to each other

$$\frac{\mathcal{H}_{EWP}^{(\Delta S=1)}(\overline{15})}{\mathcal{H}_{Tree}^{(\Delta S=1)}(\overline{15})} \equiv -\frac{3c_9 + c_{10}}{2c_1 + c_2} \frac{V_{tb}^* V_{ts}}{V_{ub}^* V_{us}}$$

$$\equiv -(0.65 \pm 0.15) e^{-i\gamma} \equiv -\delta_{EW} e^{-i\gamma}$$

$\Rightarrow$  Hadronic matrix elements are related and involve a common strong phase

**Example:**  $B^+ \rightarrow K^0 \pi^+, K^+ \pi^0$

Where does only  $\overline{15}$  contribute?  $B \rightarrow (K\pi)_{I=3/2}$

$$A(B^+ \rightarrow K^0 \pi^+) + \sqrt{2}A(B^+ \rightarrow K^+ \pi^0) \equiv \langle \overline{15} \rangle$$

$$-(T + C + P_{EW} + P_{EW}^c) \equiv -(T + C)(1 - \delta_{EW} e^{-i\gamma})$$

Tree and EWP are subdominant

Dominant penguin  $\in \overline{3}$  are equal ( $\Delta I \equiv 0$ ) in

$$A(B^+ \rightarrow K^0 \pi^+) \text{ and } -\sqrt{2}A(B^+ \rightarrow K^+ \pi^0)$$

$$R_*^{-1} \equiv \frac{2\Gamma(B^\pm \rightarrow K^\pm \pi^0)}{\Gamma(B^\pm \rightarrow K^0 \pi^\pm)}$$

$$\frac{|R_*^{-1} - 1|}{2\epsilon} \leq |\cos \gamma - \delta_{EW}| \quad (\delta_{EW} \equiv 0.65 \pm \dots)$$

$$\epsilon \equiv \sqrt{2} \frac{V_{us} f_K |A(B^+ \rightarrow \pi^0 \pi^+)|}{V_{ud} f_\pi |A(B^+ \rightarrow K^0 \pi^+)|} \equiv 0.20 \pm 0.05$$

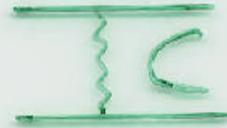
$$[\text{also } A(B^+ \rightarrow \pi^+ \pi^0) \equiv \langle |\overline{15}| \rangle \propto T + C]$$

$$(R_*^{-1})_{ave} \equiv 1.45 \pm 0.46 \neq 1 \text{ ???} \Rightarrow \gamma \neq \sim 50^\circ$$

## Final state rescattering

No Tree amplitude:

- $A(B^0 \rightarrow K^+K^-) \equiv PA + E$



Is  $E$  small  $\propto f_B/m_B$ ? How large is rescattering

$$B^0 \rightarrow \pi^+\pi^-, \rho^+\rho^- \rightarrow K^+K^- ?$$

No rescattering:  $\mathcal{B}(B^0 \rightarrow K^+K^-) \equiv n \times 10^{-8}$

Present limit  $\mathcal{B}(B^0 \rightarrow K^+K^-) < 1.9 \times 10^{-6}$   
needs improvement

- $A(B^+ \rightarrow K^0\pi^+) \equiv (P - \frac{1}{3}P_{EW}^c) + A$



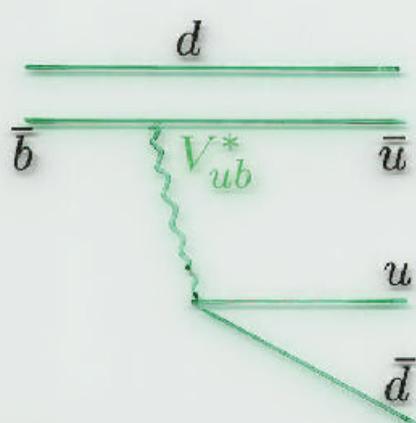
Is  $A$  small  $\propto f_B/m_B$ ? How large is rescattering

$$B^+ \rightarrow K^+\pi^0, K^{*+}\rho^0 \rightarrow K^0\pi^+ ?$$

$$A(B^+ \rightarrow \bar{K}^0K^+) \equiv -(P - \frac{1}{3}P_{EW}^c)\tan\theta_c + A\tan^{-1}\theta_c$$

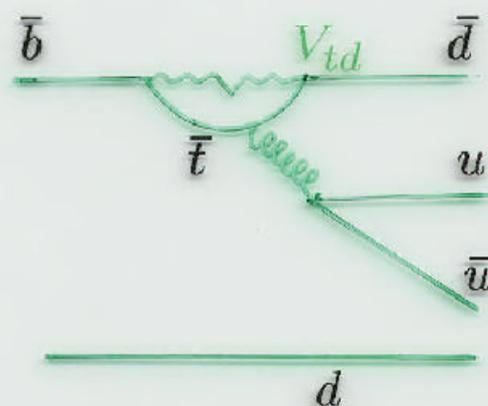
enhancement by  $\tan^{-2}\theta_c \approx 20$  of  $A/P$  ratio in  
 $B^+ \rightarrow \bar{K}^0K^+$  relative to  $B^+ \rightarrow K^0\pi^+$  can be  
measured

•  $B^0(t) \rightarrow \pi^+\pi^-$ : Penguin Pollution



$$V_{ub}^* V_{ud} \langle T \rangle$$

phase  $\gamma$



$$V_{tb}^* V_{td} \frac{\alpha_s}{12\pi} \ln \frac{m_t^2}{m_b^2} \langle P \rangle$$

phase  $= \beta$

CLEO '98:  $K\pi/\pi\pi$ :  $\frac{\text{Penguin}}{\text{Tree}} \sim 0.3 \pm 0.1$

$$A(t) \equiv a_{\text{dir}} \cos(\Delta mt) + \sqrt{1 - a_{\text{dir}}^2} \sin 2(\alpha + \theta) \sin(\Delta mt)$$

direct CP violation

penguin pollution

$$a_{\text{dir}} \sim \frac{\text{Penguin}}{\text{Tree}} \sin \delta$$

$$\theta \sim \frac{\text{Penguin}}{\text{Tree}} \cos \delta$$

$$\delta \equiv \text{Arg}(P/T)$$

- **solution 1: Isospin symmetry**

$$B^0(t) \rightarrow \pi^+ \pi^-, B^+ \rightarrow \pi^+ \pi^0, B^0 \rightarrow \pi^0 \pi^0, C.C.$$

$$- P \text{ is pure } \Delta I = 1/2 \Rightarrow I(\pi\pi) = 0$$

$$- T \propto P^{EW} \Rightarrow I(\pi\pi) = 2$$

### Isospin triangle resolves penguin pollution

Difficulty:  $B^0 \rightarrow \pi^0 \pi^0$  is color-suppressed  $BR < 10^{-6}$  (if little rescattering  $B \rightarrow D^+ D^-, \pi^+ \pi^- \rightarrow \pi^0 \pi^0$ )

$$\sin(\delta\alpha) \leq \sqrt{\frac{\mathcal{B}(B \rightarrow \pi^0 \pi^0)}{\mathcal{B}(B^\pm \rightarrow \pi^\pm \pi^0)}} \quad (\text{Grossman, Quinn})$$

- **solution 2: Calculate  $|P/T|$  ( $\delta = \text{Arg}(P/T)$ )**

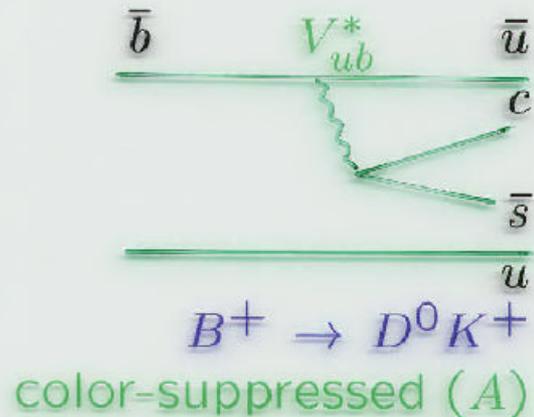
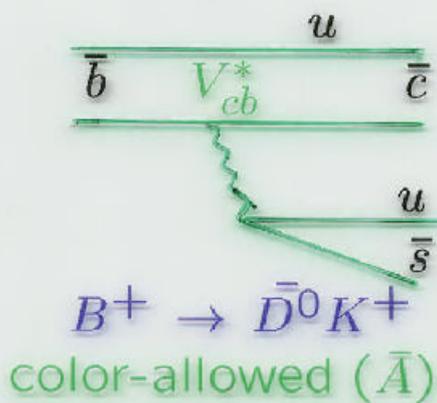
$$- a_{\text{dir}}, a_{\text{mix}} \equiv \text{functions of } (|P/T|, \delta, \beta, \gamma)$$

- after measurement of  $\beta$ , measurement of  $a_{\text{dir}}, a_{\text{mix}}$  determines  $\delta$  and  $\gamma$

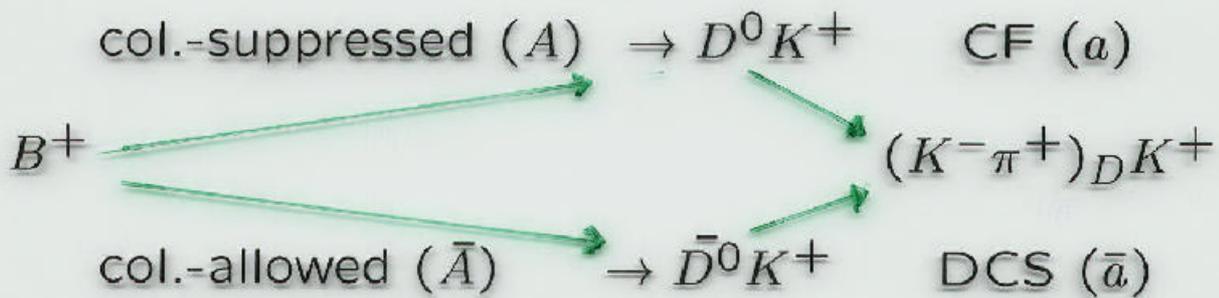
- **New physics**

(D. London)

- $B^\pm \rightarrow DK^\pm$  measures  $\sin \gamma$  precisely



- (1)  $D^0$  flavor modes:  $D^0 \rightarrow K^- \pi^+$



$$\Gamma \equiv \Gamma(B^+ \rightarrow (K^- \pi^+)_{D} K^+) \equiv |Aa|^2 + |\bar{A}\bar{a}|^2 + 2|A\bar{A}a\bar{a}| \cos(\delta_{DK} + \Delta_{K\pi} + \gamma)$$

$$\bar{\Gamma} \equiv \Gamma(B^- \rightarrow (K^+ \pi^-)_{D} K^-) \equiv \Gamma(\gamma \rightarrow -\gamma)$$

$$\frac{\bar{\Gamma} - \Gamma}{\bar{\Gamma} + \Gamma} \equiv \frac{2|\bar{A}\bar{a}Aa| \sin(\delta + \Delta) \sin \gamma}{|Aa|^2 + |\bar{A}\bar{a}|^2 + 2|A\bar{A}a\bar{a}| \cos(\delta + \Delta) \cos \gamma}$$

maximal for  $|\bar{A}\bar{a}/Aa| \equiv 1, \delta + \Delta \equiv \pi/2$

measured  $\bar{A}, a, \bar{a}$ :

$$|\bar{A}|^2 \propto \mathcal{B}(B^+ \rightarrow \bar{D}^0 K^+) \equiv (2.9 \pm 0.8) \times 10^{-4}$$

$$|a|^2 \propto \mathcal{B}(D^0 \rightarrow K^- \pi^+) \equiv (3.83 \pm 0.09) \times 10^{-2}$$

$$|\bar{a}|^2 \propto \mathcal{B}(\bar{D}^0 \rightarrow K^- \pi^+) \equiv (1.46 \pm 0.30) \times 10^{-4}$$

$$\frac{|\bar{a}|}{|a|} \equiv 0.062 \pm 0.007 \equiv (1.21 \pm 0.13) \tan^2 \theta_c$$

$$\equiv \tan^2 \theta_c \quad \text{U-spin}$$

not measured:  $A, (\delta_{DK} + \Delta_{K\pi}), \gamma$

$$\frac{|A|}{|\bar{A}|} \sim \frac{|V_{ub}^* V_{cs}| |a_2|}{|V_{cb}^* V_{us}| |a_1|} \sim 0.1 \quad \Rightarrow \quad \frac{|\bar{A}\bar{a}|}{|Aa|} \sim 0.6$$

$\delta_{DK} \equiv ?$  probably small, as in  $B \rightarrow D\pi$

$$\Delta_{K\pi} \equiv \text{Arg} \frac{A(D^0 \rightarrow K^- \pi^+)}{A(\bar{D}^0 \rightarrow K^- \pi^+)} \equiv ?$$

$$\equiv 0 \quad \text{U-spin, estimate } < 20^\circ$$

$\Rightarrow$  sizable asymmetry cannot be excluded

$\Gamma(B^\pm \rightarrow (K\pi)_D K^\pm) \equiv$  functions of  $A, (\delta + \Delta), \gamma$

determination of  $\gamma$  requires another DCS mode

$$\bar{D}^0 \rightarrow K^- \pi^+ \pi^0 \quad (K^* \pi, K\rho) \quad (\text{A. Smith})$$

(2)  $D^0$  CP modes:  $D_{\pm}^0 \rightarrow K^+ K^-$

$$D_{+,-}^0 \equiv \frac{1}{\sqrt{2}}(D^0 \pm \bar{D}^0)$$

$$\sqrt{2}A(B^+ \rightarrow D_{\pm}^0 K^+) \equiv \pm |\bar{A}| + |A| \exp[i(\delta + \gamma)] \quad .$$

$$R_{\pm} \equiv \frac{2[\Gamma(B^+ \rightarrow D_{\pm}^0 K^+) + C.C.]}{\Gamma(B^+ \rightarrow \bar{D}^0 K^+) + C.C.}$$

$$A_{\pm} \equiv \frac{\Gamma(B^+ \rightarrow D_{\pm}^0 K^+) - C.C.}{\Gamma(B^+ \rightarrow \bar{D}^0 K^+) + C.C.} \quad .$$

$$R_{\pm} \equiv 1 + |A/\bar{A}|^2 \pm 2|A/\bar{A}| \cos \delta \cos \gamma$$

$$A_{-} \equiv -A_{+} \equiv |A/\bar{A}| \sin \delta \sin \gamma$$

Three equations for  $|A/\bar{A}|$ ,  $\delta$ ,  $\gamma$

$$1 + |A/\bar{A}|^2 \equiv (R_{+} + R_{-})/2$$

cannot measure  $|A/\bar{A}| \sim 0.1$  precisely

$$\sin^2 \gamma \leq R_{\pm} \Rightarrow \text{new constraint on } \gamma$$

$$\text{e.g.: } \gamma \equiv 40^\circ, \frac{|A|}{|\bar{A}|} \equiv 0.1, \delta \equiv 0 \Rightarrow R_{-} \equiv 0.85$$

$$10^8 B^+ B^- \text{ pairs: } \delta(R_{-}) \equiv 0.05$$

$$\Rightarrow 73^\circ < \gamma < 107^\circ \text{ excluded at 90\% c.l.}$$

- combine  $D^0$  CP modes and  $D^0$  flavor modes

(Atwood, Dunietz, Soni)

(Soffer)

- measure  $\sin(2\beta \pm \gamma)$  in  $B^0(t) \rightarrow D^{*0} K^{*0}$  through time and angular dependence involving interference of helicity amplitudes

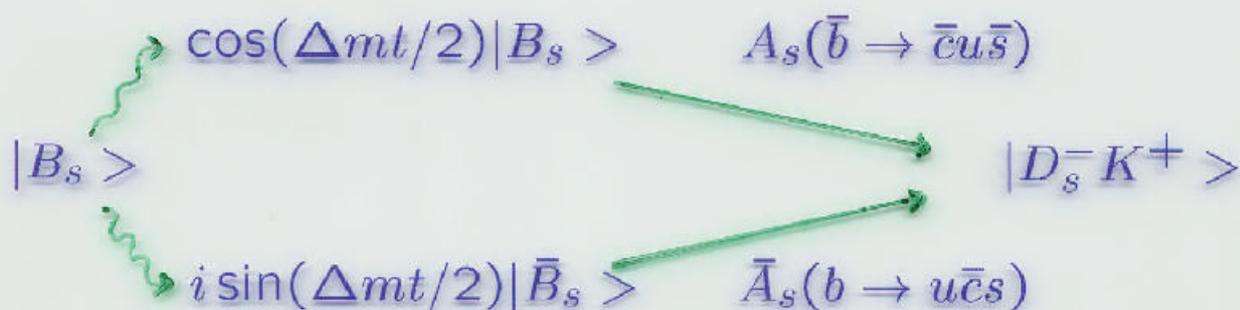
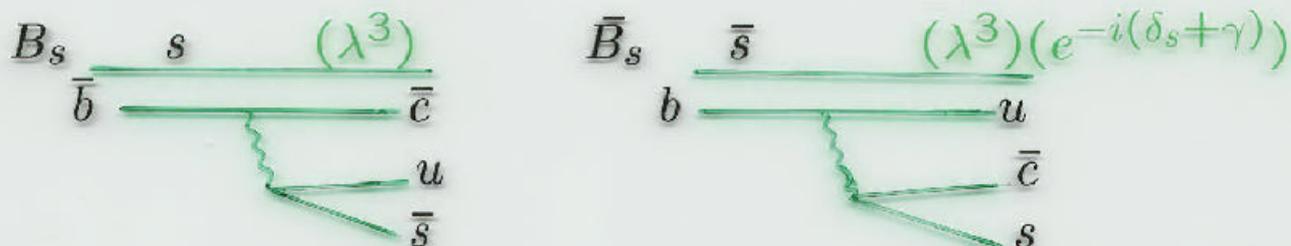
(London, Sinha<sup>2</sup>)

- $\mathcal{B}(B^+ \rightarrow \bar{D}^0 K^+) \mathcal{B}(\bar{D}^0 \rightarrow K^- \pi^+) \equiv (4.2 \pm 1.4) \times 10^{-8} \Rightarrow$  **method requires  $> \mathcal{O}(10^8) B\bar{B}$  pairs**

(D. Atwood)

- $B_s(t) \rightarrow D_s K$  **measures**  $\sin \gamma$

(1) flavor tagged  $B_s$  time-dependence (*Dunietz et al*)



$$\Gamma(B_s(t) \rightarrow D_s^- K^+) \equiv e^{-\Gamma_s t} \left[ |A_s|^2 \cos^2\left(\frac{\Delta m t}{2}\right) + |\bar{A}_s|^2 \sin^2\left(\frac{\Delta m t}{2}\right) - |A_s \bar{A}_s| \sin(\delta_s + \gamma) \sin(\Delta m t) \right]$$

$$(a) \Gamma(B_s(t) \rightarrow D_s^- K^+) \equiv f(|A_s|, |\bar{A}_s|, \delta_s + \gamma)$$

$$(b) \Gamma(\bar{B}_s(t) \rightarrow D_s^+ K^-) \equiv f(\gamma \rightarrow -\gamma)$$

$$(c) \Gamma(\bar{B}_s(t) \rightarrow D_s^- K^+) \equiv f(|A_s| \leftrightarrow |\bar{A}_s|)$$

$$(d) \Gamma(B_s(t) \rightarrow D_s^+ K^-) \equiv f(|A_s| \leftrightarrow |\bar{A}_s|, \gamma \rightarrow -\gamma)$$

this determines  $\gamma$

(2) Untagged  $B_s(t)$  including  $\frac{\Delta\Gamma_s}{\Gamma_s} \sim 10\%$

$$\Gamma(B_s(t) \rightarrow D_s^- K^+) + \Gamma(\bar{B}_s(t) \rightarrow D_s^- K^+) \equiv \frac{1}{2}(|A_s|^2 + |\bar{A}_s|^2)(e^{-\Gamma_L t} + e^{-\Gamma_H t}) + |A_s \bar{A}_s| \cos(\delta_s + \gamma)(e^{-\Gamma_L t} - e^{-\Gamma_H t})$$

$$\Gamma(B_s(t) \rightarrow D_s^+ K^-) + \Gamma(\bar{B}_s(t) \rightarrow D_s^+ K^-) : \gamma \rightarrow -\gamma$$

2 rates determine  $\gamma$  for known  $|A_s|$

( $|A_s|$  obtained from tagged  $B_s$  or factorization)

(3) CP-tagged  $B_s$  time-integrated (Falk)

$$B_s^{CP+} \equiv (B_s + \bar{B}_s)/\sqrt{2} \quad B_s^{CP+} \rightarrow D_s^+ D_s^-$$

$$A(B_s^{CP+} \rightarrow D_s^- K^+) \equiv [ |A_s| + |\bar{A}_s| e^{-i(\delta_s + \gamma)} ] / \sqrt{2}$$

$$A(B_s^{CP+} \rightarrow D_s^+ K^-) \equiv [ |A_s| + |\bar{A}_s| e^{-i(\delta_s - \gamma)} ] / \sqrt{2}$$

two triangles determine  $\gamma$

requires independent measurement of  $|A_s|$ ,  $|\bar{A}_s|$

CP-tagged  $B_s^{CP+} \rightarrow D_s^+ K^-$  can be measured at a very high luminosity  $e^+e^-$  B factory on  $\Upsilon(5S)$

## conclusions

- Measure crudely many CP asymmetries

$$\Delta(B^0 \rightarrow K^+ \pi^-) \equiv -\Delta(B_s \rightarrow \pi^+ K^-)$$

$$\Delta(B^0 \rightarrow \pi^+ \pi^-) \equiv -\Delta(B_s \rightarrow K^+ K^-)$$

relative signs test CKM (U-spin)

- Measure precisely certain charge-averaged rates

$$\frac{\Gamma(B^\pm \rightarrow K^\pm \pi^0)}{\Gamma(B^\pm \rightarrow K \pi^\pm)}, \quad \frac{\Gamma(B^\pm \rightarrow D_{CP}^0 K^\pm)}{\Gamma(B^\pm \rightarrow D_{FL}^0 K^\pm)}$$

constrain  $\phi_3 \equiv \gamma$

- Measure precisely certain CP asymmetries

$$B \rightarrow \pi\pi, K\pi, DK, B_s \rightarrow K\pi, K\bar{K}, D_s K$$

determine  $\phi_3 \equiv \gamma$  precisely

- Improve theory of hadronic B decays (*Brodsky, Li, Keum, Sanda, Silvestrini*)

calculate  $P/T$ , SU(3) breaking, rescattering

questions can be studied experimentally