

# Charm Physics and the Poor Sleeper's Impatience

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## Disclaimer:

The fascination of charm decays  
cannot match that of beauty decays  
anymore than Botticelli can match Michelangelo!  
yet:

- test bed for QCD technologies
- charm transitions are a **unique** portal for obtaining a **novel** access to the **flavour problem** with the **experimental situation** being a priori favourable!

## topics

- QCD technologies
- On charm's promise of New Physics
- Appeal for a comprehensive New Phenomenology
  - $D^0 - \bar{D}^0$  oscillations
  - ~~CP~~

# I. Theoretical Technologies for QCD

s	$m_s < \Lambda$	chiral pert. th.		
c	$m_c > \Lambda$	$1/m_c$ expans. (?)	↑	↓
b	$m_b \gg \Lambda$	$1/m_b$ expans. !		↑
t	$\Gamma_t \sim O(\Lambda)$	perturb. dynamics	$1/m_Q$ exp.	lattice QCD

- chiral pert. theory
  - lattice QCD
  - $1/m_Q$  expansions
  - quark models (properly used)
- ground prepared for fruitful feedback

## I.1 Heavy quark expansions

$$\gamma(E) \equiv d\Gamma/d\dots \equiv F \sum_i c_i (\alpha_S) (\Lambda_i/m_Q)^i$$

tools

- effective Lagrangians

$$\mathcal{L}(\Lambda_{UV}) \rightarrow \mathcal{L}(m_Q) \rightarrow \mathcal{L}(\mu)$$

$$\Lambda_{UV} \gg m_Q \gg \mu$$

- operator product expansion (OPE)
- sum rules

$$\int_0^\infty dE w(E) \gamma(E)|_{\text{quarks}} \equiv \int_0^\infty dE w(E) \gamma(E)|_{\text{hadrons}}$$

applied to

- ➔ **inclusives**: lifetimes, SL BR's, lepton **spectra**  
report card -- later

- ➔ **exclusives**:

- SL form factors:

HQET -- no **absolute** predictions

- NL two-body modes

models still only tool available

(new framework for  $B \rightarrow M_1 M_2$  hard to justify here  
-- but should be tried anyhow!)

## I.2 Quark-hadron duality (duality)

$$\langle d\sigma(\text{quark\&gluon d.o.f.}) \rangle \equiv \langle d\sigma(\text{hadr.d.o.f.}) \rangle$$

 duality

no complete theory yet for duality and its limitations -- but we have moved beyond the folkloric stage in the last few years

we understand physical origins

- hadronic thresholds
- 'distant' cuts
- $1/m_c$  expansions

we have identified mathematical portals

Euclidean  $\exp\{-m_Q/\Lambda\}$    
Minkowskian  $\sin(m_Q/\Lambda)$

duality -- a very natural concept:

hadronic final state formed in 2 steps

① hard process in femto universe --

time scale  $\sim 1/m_Q$

② hadronization soft --

$1/\Lambda$  in restframe of Q  $\Rightarrow \sim m_Q/\Lambda^2$  in c.m. frame

$$m_Q/\Lambda^2 \gg 1/m_Q$$

gross features determined by 1st step

➔ duality exact at asymptotic scales

💣 corrections at finite scales!

illustration by qm model with potential  $V(\mathbf{x})$ :

○ local properties of  $V \rightarrow$  integrated rates \*

○ asymptotic “ of  $V \rightarrow$  specifics of final state

\* unless there are singularities at finite distances

probes:

- OPE  $\longrightarrow$  indirect lessons
  - dispers. relat.  $\rightarrow$  sum rules  $\nearrow$
  - resonance models  $\longrightarrow$  direct lessons
  - t'Hooft model  $\nearrow$
- ( $\equiv$  QCD in 1+1 dimensions with  $N_C \rightarrow \infty$ )

- *over-constraining measurements*

*final arbiter !*

## general findings:

- ➔ duality cannot be exact at finite scales
- ➔ limitations to duality will depend on the process:  $\sim \sin(m_Q/\Lambda)/m_Q^k$ ,  $k > 1$
- ➔ fundamental question:  
is there an OPE -- or not!
- ➔ duality violations larger in NL than SL decays, but no fundamental difference!
- ➔ difference between local and other versions of duality quantitative rather than qualitative

## one particular and obvious problem in charm sector:

expansion parameter  $\Lambda_1/m_c \leq 1$

since  $m_c(m_c) = 1.25 \pm 0.1 \text{ GeV}$

- ➔ at best: sizeable uncertainties
- at worst: duality inoperative at charm scale

## I.3 Lattice QCD

originally introduced

- to prove confinement and
- bring spectroscopy under theoret. control

now making major contributions to heavy flavour physics... with **partially unquenched** results

→ decay constants  $255 \pm 30$

$$f(D_s) \equiv \begin{cases} \cancel{240 \pm 4 \pm 24}, \cancel{275 \pm 20} \text{ MeV, lattice } n_f=2 \\ 269 \pm 22 \text{ MeV, exp. WA,} \end{cases}$$

↙ probably underestimate!

expectation for the future:

$$\Delta f(D^+), \Delta f(D_s) \sim \begin{cases} 6 - 10 \% , \text{ BABAR/BELLE} \\ \text{using CLEO method} \\ 2 - 3 \% , \tau\text{-charm factory} \end{cases}$$

analysis with **full unquenching** that treats **charm quarks dynamically** not utopian

➔ semileptonic form factors

$$D \rightarrow l \nu P$$

in the future:

analysis with full unquenching that treats charm quarks dynamically required and not utopian

my expectation:

charm decays provide rich lab for quantitative tests of lattice QCD

## I.4 Synergies

natural feedback between the two technologies

- both defined in Euclidean space
- both “mature”
- similar as well as different expansion parameters
- lattice QCD provides input to  $1/m_Q$  expansion

high sensitivity → high accuracy probe for NP!

• charm scale ~ bridge between

$1/m_Q$  expansion and lattice QCD

## II. Present Profile of Weak Dynamics of Charm

### ○ lifetimes

### hierarchy

$$\tau(D^+) > \tau(D^0) \sim \tau(D_s^+) \geq \tau(\Xi_c^+) > \tau(\Lambda_c^+) > \tau(\Xi_c^0) > \tau(\Omega_c^+)$$

	$1/m_c$ expect.	comments on expect.	data * PDG '00 ** Osaka '00
$\tau(D^+)/\tau(D^0)$	$\sim 2$	PI in $\tau(D^+)$	$2.54 \pm 0.03$ **
$\tau(D_s)/\tau(D^0)$	1.0 – 1.07 0.9 – 1.3 $1.08 \pm 0.04$	without WA with WA QCD SR	$1.125 \pm 0.042$ PDG '98 $1.18 \pm 0.02$ **
$\tau(\Lambda_c^+)/\tau(D^0)$	$\sim 0.5$	Quark model matrix elem.	$0.50 \pm 0.03$ *
$\tau(\Xi_c^+)/\tau(\Lambda_c)$	$\sim 1.3$	**	$1.60 \pm 0.30$ * 2.14, FOCUS
$\tau(\Xi_c^+)/\tau(\Xi_c^0)$	$\sim 2.8$	**	$3.37 \pm 0.91$ * 3.56, FOCUS
$\tau(\Xi_c^+)/\tau(\Omega_c)$	$\sim 4$	**	$5.2 \pm 1.8$ *

### ➔ semilept. BR's

$$BR_{SL}(D^0) = 6.7\% \quad \text{vs.} \quad \sim 8\%$$

$$BR_{SL}(D^0) = 17.2\% \quad \text{vs.} \quad \sim 16\%$$

$$\text{OPE term} \sim 1/m_c^2$$

## Score card:

- predictions better than could have been counted on.
- $1/m_Q$  expansions provide an after the fact rationale for most phenomenological concepts like **PI**, **WA**, **WS** etc. which are  $\sim O(1/m_Q^3)$
- they are more definite about those concepts; e.g., WA has to be **non**leading, though still significant.
- predictions for baryon lifetimes are based on QM calculations of expectation values
- need more precise data on  $\Xi_c^0$  and  $\Xi_c^-$  lifetimes
- the **absolute**  $D^0$  &  $D^+$  SL BR's are understood now due to an  $O(1/m_Q^2)$  effect
- ratios of SL BR's of baryons do **not** reflect their lifetime ratios

○ *Cabibbo Hierarchy*

$H_c \rightarrow l\nu$  [ $S \equiv -1, 0$ ]      observed

$H_c \rightarrow$  [ $S \equiv -1, 0, +1$ ]      “

○  $V(cs), V(cd)$

*imposing 3-family unitarity*

$$|V(cs)| \equiv 0.9742 \pm 0.0008$$

$$|V(cd)| \equiv 0.222 \pm 0.003$$

*without imposing 3-family unitarity*

$$|V(cs)| \equiv 0.880 \pm 0.096 \text{ from } D \rightarrow l\nu K$$

$$\nu N \rightarrow l^+ l^- X$$

$$|V(cd)| \equiv 0.226 \pm 0.007 \text{ from } \nu N \rightarrow l^+ l^- X$$

$$D \rightarrow l\nu K$$

**new OPAL analysis of  $W \rightarrow H_c X$**

$$|V(cs)| \equiv 0.969 \pm 0.058$$

loop processes ...

general expectations:

- ☞ slow  $D^0 - \bar{D}^0$  oscillations
- ☞ tiny CP asymmetries
- ☞ extremely rare decays
- ➔ ~ zero background search for New Physics

☐  $D^0 - \bar{D}^0$  oscillations

controlled by two quantities

$$x_D \equiv \frac{\Delta m_D}{\Gamma_D} \quad y_D \equiv \frac{\Delta \Gamma_D}{2\Gamma_D}$$

$$\Gamma_D = \frac{\Gamma(D^0 \rightarrow l^- X)}{\Gamma(D^0 \rightarrow l^+ X)} \approx \frac{x_D^2 + y_D^2}{2} \quad \text{for } x_D, y_D \ll 1$$

- $x_D$  and  $y_D$  Cabibbo suppressed
- $x_D \equiv 0 \equiv y_D$  in the  $SU(3)_{\text{Fl}}$  limit
- ➔  $x_D, y_D < 0.05$

a conservative estimate

$$x_D, y_D \sim O(0.01)$$

data

$$y'_D \equiv -2.5^{+1.4}_{-1.6} \pm 0.3 \%$$

CLEO,  $D \rightarrow K\pi$

$$y_D \equiv 0.8 \pm 2.9 \pm 1.0 \%$$

E 791,  $D \rightarrow l\nu K$

$$y_D \equiv 1.0^{+3.8}_{-3.5} {}^{+1.1}_{-2.1} \%$$

BELLE

$$y_D \equiv 3.42 \pm 1.39 \pm 0.74 \%$$

FOCUS,  $D \rightarrow KK$

vs.  $D \rightarrow K\pi$

$$y_D \equiv -1.1 \pm 2.5 \pm 1.4 \%$$

CLEO,  $D \rightarrow KK, \pi\pi$

vs.  $D \rightarrow K\pi$

i.e., data consistent with zero -- on the % level

❑ ~~CP~~

time integrated partial widths

possible with SM & KM in singly Cabibbo suppressed modes

benchmark estimate:

asymmetry  $\sim O(\lambda^4) \sim O(0.001)$

data, WA Osaka '00

$$A_{CP}(D^0 \rightarrow K^+K^-) = 0.5 \pm 1.6 \%$$

$$A_{CP}(D^0 \rightarrow \pi^+\pi^-) = 2.2 \pm 2.6 \%$$

$$A_{CP}(D^\pm \rightarrow K^\pm K^- \pi^+) = 0.2 \pm 1.1 \%$$

CLEO '01

$$A_{CP}(D^0 \rightarrow K^+K^-) = 0.05 \pm 2.18 \pm 0.84 \%$$

$$A_{CP}(D^0 \rightarrow \pi^+\pi^-) = 2.0 \pm 3.2 \pm 0.8 \%$$

$$A_{CP}(D^0 \rightarrow K_S \pi^0) = 0.1 \pm 1.3 \%$$

$$A_{CP}(D^0 \rightarrow \pi^0 \pi^0) = 0.1 \pm 4.8 \%$$

i.e., data consistent with zero -- on the % level

## Interlude -- Info on Charm as Engineering Input

- unfinished business
  - ▢ absolute BR's, in particular for  $D_s$  and charm baryons
  - ▢  $\Xi_c$  and  $\Omega_c$  lifetimes
  - ▢ SL BR's of charm baryons
  - ▢ more precise data on  $D_s \rightarrow l \nu$  and measurements of  $D^+ \rightarrow l \nu$
  - ▢ post-MARK III data on lepton spectra in inclusive SL charm decays

all important engineering inputs to beauty studies and some provide interesting lessons on QCD, but ...

is it more than dotting the "i"'s and crossing the "t"'s?

## III. Charm Decays -- Novel Portals to New Physics

### III.1 General remarks

SM incomplete!

where to look for New Physics?

- charm only **up-type** quark allowing full range of **indirect** searches for New Physics
  - $D^0 - \bar{D}^0$  oscillations      **no**  $T^0 - \bar{T}^0$  oscill.  
[**no** top hadroniz.]
  - ~~CP~~ with  $D^0 - \bar{D}^0$  oscill.      **cannot** occur
  - direct ~~CP~~ in excl. modes with decent BR's      tiny BR's,  
lost coherence
  - charm decays proceed in resonance region  
⇒ FSI of great vitality:  
optimal for getting signals  
(not for interpreting them)

- practical advantages and opportunities
  - ▣ large rates
  - ▣ long lifetimes
  - ▣  $D^* \rightarrow D\pi$  flavour tagging

basic contention:

charm transitions are a **unique** portal for obtaining  
a **novel** access to the **flavour problem**  
with the **experimental situation** being a priori  
**favourable!**

## III.2 $D^0 - \bar{D}^0$ oscillations

a tough challenge for theoret. technologies  
as well!

○

$$x_D = \frac{\Delta m_D}{\Gamma_D}$$

$$y_D = \frac{\Delta \Gamma_D}{2\Gamma_D}$$

general bound:  $x_D, y_D < 0.05$

a conservative estimate :  $x_D, y_D \sim O(0.01)$

$$y'_D \equiv -2.5^{+1.4}_{-1.6} \pm 0.3 \%$$

CLEO,  $D \rightarrow K\pi$

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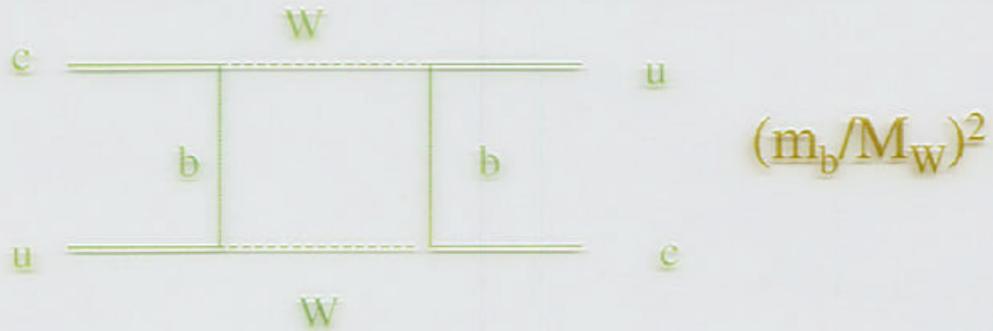
vs.  $D \rightarrow K\pi$

the game has just begun!

considerable previous literature

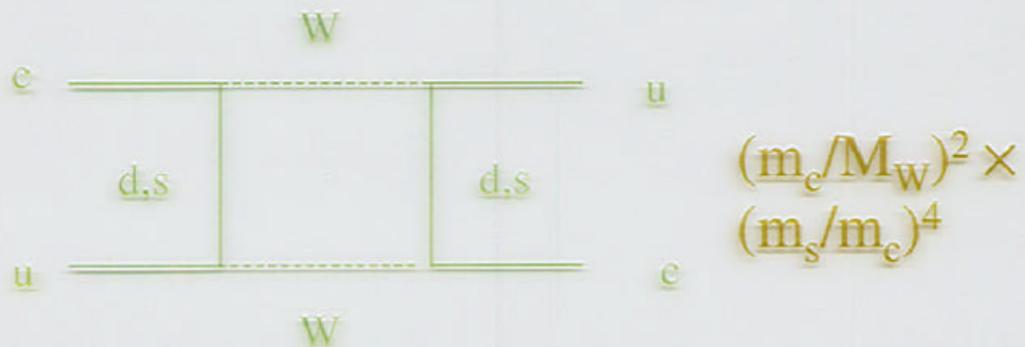
○ quark box diagram

□ local operator:  $\sim m_b > M_D$



$$x_D(\text{bb}) \sim \text{few} \times 10^{-7}$$

□ short distance operator:  $\sim m_c > \bar{\Lambda}$



$$x_D(\text{box}) \sim \text{few} \times 10^{-5}$$

○ long distance contributions:

use various schemes to describe selected hadronic states:

➔  $x_D(\text{LD}), y_D(\text{LD}) \sim 10^{-4} - 10^{-3}$

systematic analysis based on  $O_{\text{perator}}$   $P_{\text{roduct}}$   $E_{\text{xpansion}}$

$$T(\Delta C=2, \omega) \equiv \int d^4x e^{-i\omega t} \langle D | \{ \mathcal{L}(x) \mathcal{L}(0) \}_T | D \rangle / 4M_D$$

$$4 T(\Delta C=2, \omega) \equiv A(\omega) + A(-\omega)$$

$$A(\omega) \equiv -\Delta \overline{M}_D(\omega) + i \Delta \overline{\Gamma}_D(\omega) / 2$$

$$\Delta M_D \equiv \Delta \overline{M}_D(0), \Delta \Gamma_D \equiv \Delta \overline{\Gamma}_D(\omega)$$

$$\Delta M_D \equiv (1/2\pi) \text{P.V.} \int d\omega [\Delta \overline{\Gamma}_D(\omega) / \omega]$$

→ expansion in powers of  $1/m_c, m_s, KM$



GIM suppression  $(m_s/m_c)^4$  of usual quark box diagram **untypically severe!**

∃ contributions from **higher-dimensional operators** with a **very gentle GIM factor**

$\sim m_s / \mu_{\text{had}}$  ... due to condensates in the OPE